"材料与能源前沿科学:激发态和动力学"培训班, 2024.05.13

关于原子核的量子效应的模拟方法的介绍

李新征

北京大学物理学院

感谢高老师的邀请!



●什么是原子核的量子效应?

●模拟原子核量子效应的计算方法有哪些,它们的优缺点是什么?

●几个例子,来感受相关研究。

Take-home Message.

第一部分:什么是原子核的量子效应?

●物理学的核心任务是思维范式的建立!



●出发点:上面那句话。

Apply to: 牛顿力学、热力学、电动力学、量子力学、统计力学、相对论



学完之后,我们脑子里面无非是有了一个理论框

架,去描述相关问题(客观、有预测能力)

第一部分:什么是原子核的量子效应?

历史上,我们凝聚态物理的发展,一模一样走的就是这条路!



第一部分:什么是原子核的量子效应?

上世纪40年代开始,高潮是50、60年代,之后依然延展



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第一部分:什么是原子核的量子效应?

上世纪50年代开始,70年代喊出来



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1958 Anderson Localization Mott (1949)-Hubbard (1963) Insulator

Emergent Phenomenon

第一部分:什么是原子核的量子效应?



第一部分:什么是原子核的量子效应?



第一部分:什么是原子核的量子效应?

● 这个情况,同样apply to凝聚态物理

Coupled nuclear & electronic DOFs



Franck-Condon principle (1926):



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Clear: Quantum states of vibration contribute to the spectrum. Unclear: how should the quantum state with nuclei be rigorously expressed?

第一部分:什么是原子核的量子效应?

● 这个情况,同样apply to凝聚态物理

Die Gültigkeitsgrenze der Theorie der idealen Kristalle und ihre Überwindung.

Von Max Born, F. R. S., Edinburgh.

Einleitung.

Die Theorie der idealen Kristalle, an der ich seit mehr als 40 Jahren gearbeitet habe, ist trotz mancher Erfolge weder logisch noch empirisch befriedigend. Die wichtigsten Einwände sind die folgenden:

DYNAMICAL THEOR CRYSTAL LATTI

BY

KUN HUANG

OXFORD AT THE CLARENDON PRI

	143			
OF	1957			
Or	10.9	VIII		
ES	10.0	ELIMINATION OF THE ELECTRONIC MOTIO	N (p. 172)	
		In the text (see § 14) and in Appendix VII, this have of molon for the mach are obtained by a systematic expansion in powers of the parameter $s = (m/Q)$ (14.5), and it is shown that the adiabatic approximation, in which the electronic motion is calculated as if the molel were rest, is valid ip to terms of the form order in s. The averaged electronic energy (eigenvalue) in a given stars plays, to this approximation, the part of the poleration of the molel. It have, however, been found, in particular by studying molecular transitions that the adiabatic models has a write application than predicted by this theory. There exists in fact another motion which contains this prestriat possible with the value of the determine since commission. Use a sightly different quantity. This method has the further advantage that it leads to a system of simultaneous equations for all electronic states which represent the coupling of electronic and equations for all electronic states which represent the coupling of electronic and equations for all electronic states which represent the coupling of electronic and the states of the intervence in the present the coupling of electronic and equations for all electronic states which represent the coupling of electronic and the states of the state		
	10.63	nuclear motion in a rigorous way. In the notation of IV 3.14 the total Hamiltonian is		
	2.5	$H = T_{\pi} \pm T_{\pi} \pm U(x X)$	OTHA	
	12.53	and that corresponding to fixed nuclei		
	283	$H^{0} = T_{y} + U(x, X),$	(VIII.2	
	1000	It is assumed, as in IV, § 14, that the latter problem is solved; in t	he equation	
	1.06	$(H^0 - \Phi_n(X))\phi_n(x, X) = 0$	(VIII.3	
	10	the functions $\Phi_{\rm s}(X)$ and $\phi_{\rm s}(x, X)$, which represent the energy as function of the electrons in the state n for a fixed nuclear configur regarded as known. The actual wave equation is	and the wave vation X , are	
	10 V G	$(H-E)\Psi(x, X) = 0;$	(VIII.4	
	1202	we try to solve it by an expansion		
	10.3	$\Psi(x, X) = \sum \phi_n(X)\phi_n(x, X).$	(VIII.5	
	184	Substituting this in (VIII.4), multiplying the result by $\phi_n^*(x, X)$ and integrating over x we obtain, with $T_X = \frac{1}{4} \sum P_n^2 M_n$, (14.1):		
	12	$(T_X+\Phi_{\mathfrak{s}}(X)-E)\phi_{\mathfrak{s}}(X)+\sum_{\mathfrak{s}'}C_{\mathfrak{s}\mathfrak{s}'}(X,P)\phi_{\mathfrak{s}'}(X)=0,$	(VIII.6	
A STREET	224	where $C_{sw} = \sum_{\mathbf{k}} \frac{1}{M_k} (A_{sw}^{cb} P_k + B_{sw}^{cb})_s$	(VIII.7,	
	10.43	and $A_{ac}^{(0)}(X) = \left(\phi_{*}^{*}(x, X) P_{*} \phi_{-}(x, X) dx \right)$		
111 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1	1000	$R^{0}(X) = 4 \int d^{*}(r X) P^{0} dr (r X) dr$	averate a	
	22	Consider the diagonal elements of these matrices. For stationar $\phi_{\mu}(x, X)$ can be chosen as real functions; then	y states the	
		$A^{\rm as}_{\rm as}(X)=-\frac{i\hbar}{2}\frac{\partial}{\partial X_k}\int \phi^{\rm a}_k(x,X)dx=0,$	(VIII.9)	

M. Born, Nachr. Akad. Wiss. Göttingen, Math.-Phys. Klasse IIa, Math.-phys.-chem. Abt., S. Art. Nr. 6, 1 (1951) (The validity limit of the theory of ideal crystals and their overcoming)



Born-Huang Expansion (1954)

第一部分:什么是原子核的量子效应?

● 这个情况,同样apply to凝聚态物理

 $\Psi(x, X, X^{0}) = \sum_{n} \psi_{n}(X, X^{0}) \varphi_{n}(x, X^{0}), \qquad (2.5) \quad \text{Born, 1951}$

 $\Psi(x,X) = \sum_{n} \psi_n(X) \phi_n(x,X).$

(VIII.5) B

Born & Huang, 1954

• A time-dependent many-body wave function of electrons and nuclei:

$$\Psi^{j}(\vec{r},\vec{R},t) = \sum_{n=1}^{c_{1}} \chi_{n}^{j}(\vec{R},t) \Phi_{n}(\vec{r},\vec{R})$$



原子核的量子效应、非绝热效应, 取舍



●什么是原子核的量子效应?

●模拟原子核量子效应的计算方法有哪些,它们的优缺点是什么?

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●几个例子,来感受相关研究。

Take-home Message.



$$\Psi^{j}(\vec{r},\vec{R},t) = \sum_{n=1}^{el} \chi_{n}^{j}(\vec{R},t) \Phi_{n}(\vec{r},\vec{R})$$

绝热层面的核量子效应(统计层面、考
虑动力学)、考虑了非绝热的核量子效
应(统计层面、考虑动力学)
方法:
波函数方法、路径积分方法、半经典方
法等。

- Physics behind: path-integral
 - Quantum mechanics: probability, propagator



Density matrix

$$\rho(x_N, x_0; 1/k_B T) = \sum_j \varphi_j(x_N) \,\varphi_j^*(x_0) e^{-E_j/k_B T} \qquad \widehat{H}(x) = -\frac{d^2}{dx^2} + V(x)$$
$$K(x_N, t_N; x_0, t_0) = \sum_j \varphi_j(x_N) \,\varphi_j^*(x_0) e^{-(i/\hbar)E_j(t_N - t_0)}$$

$$K(x_N, t_N; x_0, t_0) = \sum_{j} \varphi_j(x_N) \, \varphi_j^*(x_0) e^{-(i/\hbar)E_j(t_N - t_0)}$$



• Path-integral enters

$$\rho(x,x';k_BT) = \sqrt{\frac{2\pi\hbar}{mk_BTN}} \int_{x_0=x}^{x_N=x'} \left(\exp\left\{-\frac{1}{k_BT} \sum_{i=0}^{N} \left[\frac{m(k_BT)^2N}{2\hbar} (x_{i+1} - x_i)^2 + \frac{1}{N}V(x_i)\right] \right\} \right) \prod_{i=1}^{i=1} dx_i$$

Density matrix of a quantum system
$$\longrightarrow \quad \text{Density matrix of a classical polymer of N beads (images)}$$

Density function

$$\rho(x;k_BT) = \sqrt{\frac{2\pi\hbar}{mk_BTN}} \int_{x_0=x}^{x_{N+1}=x} \left(\exp\left\{ -\frac{1}{k_BT} \sum_{i=0}^{N-1} \left[\frac{m(k_BT)^2 N}{2\hbar} (x_{i+1} - x_i)^2 + \frac{1}{N} V(x_i) \right] \right\} \right) \prod_{i=0}^{i=0} dx_i$$

Density function of a quantum system

Density function of a polymer of N beads (images)







X.Z. Li, M. Probert, A. Alavi, and A. Michaelides, Phys. Rev. Lett. 104, 066102 (2010)

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- 分子模拟中的早期尝试:

分子动力学采样

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蒙特卡洛采样

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共同点:基于力场,无法描述化学键断裂这种有意思的现象

- 与第一性原理电子结构计算的结合:
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特点: Car-Parrinello MD (BO-MD的一个近似)

- Born-Oppenheimer MD的引入与应用:
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•动力学信息, Initial Value Representation (我 周四上午报告, 曾嘉熙)

第二部分:模拟方法、优缺点



·瞬子方法(方为,复旦大学;朱禹丞, 北京大学)



•包含量子交换的路径积分方法(杨数、
 王聪、何染尘,北京大学)





●什么是原子核的量子效应?

●模拟原子核量子效应的计算方法有哪些,它们的优缺点是什么?

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●几个例子,来感受相关研究。

Take-home Message.

Impact of quantum nuclear effects on H-bond strength?

- In 1950s, Ubbelohde effect (replace H with D) in H-bonded crystals.
- Liquids: water structure no consensus, and liquid HF is strengthened.
- Clusters: (HF)_n with n > 4, strengthened, otherwise, weakened, (H₂O)_n always weakened.

Question: is there a unified picture?

第三部分:相关例子,感受研究







X.Z Li etal. Proc. Natl. Acad. Sci. USA 108, 6369 (2011)

Isotope Substitution



Flexible monomer with anharmonic potential must be used if one want to use force-field method in PIMD simulations New Experiment (Inelastic electron)

tunneling spectroscopy, IETS):



Jing Guo, Jingtao Lü, Yexin Feng, Ji Chen, ..., Xin-Zheng Li*, Enge Wang*, Ying Jiang*, Science 352, 321 (2016)



The quantum nature of high pressure hydrogen



1). Wigner & Huntington, JCP(1935): Under high pressure (25 GPa), will H_2 become bcc solid? 2). Ashcroft, JPCM (2000): There will be a low-T liquid phase whose origin is due to QNEs.

第三部分:相关例子,感受研究

The quantum nature of high pressure hydrogen



- The quantum nature of high pressure hydrogen
 - □ How do we tackle this problem:
 - 1) Electronic structure
 - 2) Configuration space explored
 - 3) Nuclear Quantum Effects (NQEs)
 - 4) Hysteresis effect

The quantum nature of high pressure hydrogen



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The quantum nature of high pressure hydrogen



 Ji Chen, Xin-Zheng Li*, Qianfan Zhang, Matthew I. J. Probert, Chris J.
 84, 144515 (2011)

 Pickard, Richard J. Needs, Angelos Michaelides, and Enge Wang*, Nat.
 P. Cudazzo, et al., Phys. Rev. Lett. 100, 257001

 Commun. 4, 2064 (2013)
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Tunneling of hydrogen & water molecules

Low T

tunneling

Deep

In k

High T

1/T





Path-Integral based method like Instanton: Allows the determination of the chemical reaction rate using the Euclidean action of the polymer, instead of the classical free-energy.

$$k_{\text{inst}}Z_r(\beta) = (2\pi\hbar)^{-\frac{1}{2}} \left| \frac{\mathrm{d}^2 S[\mathbf{x}(\tau)]}{\mathrm{d}\tau^2} \right|^{\frac{1}{2}} Z^{\ddagger}(\beta) e^{-S[\mathbf{x}(\tau)]/\hbar}$$

Quantum tunneling can be addressed in an *ab initio* manner. 35

- Tunneling of hydrogen & water molecules
 - Experimental Technique: Field Emission Microscopy (FEM), Laser Optical Diffraction (LOD), Scanning Tunneling Microscopy (STM), and Helium Spin Echo (HeSE).

• Different transition curves has been reported, but why are they different is unclear.

 For example, on Ru(0001), a gradual transition from Arrhenius behavior to a T-independent regime has been reported. However, on Ni(100) and Cu(100), diffusion rates suddenly become T independent below a certain T, indicating a sharp classical-to-quantum transition.





Wei Fang, Jeremy O. Richardson*, Ji Chen, Xin-Zheng Li*, and Angelos Michaelides*, **Phys. Rev. Lett. 119**, 126001 (2017)

第三部分:相关例子,感受研究

Tunneling of hydrogen & water molecules

RESEARCH | REPORTS

sciencemag.org SCIENCE

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1 JANUARY 2016 • VOL 351 ISSUE 6268

Sieving hydrogen isotopes through two-dimensional crystals

M. Lozada-Hidalgo,^{1*+} S. Hu,¹⁺ O. Marshall,¹ A. Mishchenko,¹ A. N. Grigorenko,¹ R. A. W. Dryfe,² B. Radha,¹ I. V. Grigorieva,¹ A. K. Geim^{1*}

One-atom-thick crystals are impermeable to atoms and molecules, but hydrogen ions (thermal protons) penetrate through them. We show that monolayers of graphene and boron nitride can be used to separate hydrogen ion isotopes. Using electrical measurements and mass spectrometry, we found that deuterons permeate through these crystals much slower than protons, resulting in a separation factor of ≈ 10 at room temperature. The isotope effect is attributed to a difference of ≈ 60 milli–electron volts between zero-point energies of incident protons and deuterons, which translates into the equivalent difference in the activation barriers posed by two-dimensional crystals. In addition to providing insight into the proton transport mechanism, the demonstrated approach offers a competitive and scalable way for hydrogen isotope enrichment.

Yexin Feng, Ji Chen, Wei Fang, Enge Wang, Angelos Michaelides*, and Xin-Zheng Li*, J. Phys. Chem. Lett. 8, 6009 (2017)

Table 1. Calculated cNEB Barrier, ZPE Corrections (ΔE_{ZPE}) and Corrected Barrier (Barrier) for Proton Transfer Across Pristine and Hydrogenated Graphene and h-BN Sheets^{*a*}

	cNEB Barrier	$\Delta E_{ m ZPE}$	barrier
G _{pristine}	3.65	-0.26	3.39
G _{chair-H}	1.08	-0.07	1.06
G _{boat-H}	0.88	-0.12	0.76
G _{disordered-H}	0.79	-0.18	0.61





Tunneling of hydrogen & water molecules





Wei Fang, Ji Chen, Philipp Pedevilla, Xin-Zheng Li*, Jeremy
O. Richardson*, and Angelos Michaelides*, Nat. Commun.
11, 1689 (2020)



第三部分:相关例子,感受研究

Tunneling splitting of water clusters





●什么是原子核的量子效应?

●模拟原子核量子效应的计算方法有哪些,它们的优缺点是什么?

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●几个例子,来感受相关研究。

Take-home Message.

第四部分: Take-home message



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Matthias Scheffler



Enge Wang



Guangshan Tian



Part III: Quantum Nature of Hydrogen Bond

New Experiment (Inelastic tunneling spectroscopy, IETS):





Part III: Quantum Nature of Hydrogen Bond

New Experiment (Inelastic tunneling spectroscopy, IETS):





Part III: Quantum Nature of Hydrogen Bond

-> Large

200

180

New Experiment (Inelastic tunneling spectroscopy, IETS):

















Part II: Low-T metallic liquid hydrogen



Parameters in the Allen-Dynes equation

Consistent with:

J. M. McMahon and D. M. Ceperley, **Phys. Rev. B. 84**, 144515 (2011) P. Cudazzo, et al., **Phys. Rev. Lett. 100**, 257001 (2008)





A. M. J. Schaeffer et al., Phys. Rev. Lett. 109, 185702 (2012)







Y. X. Feng et al., J. Chem. Phys. 142, 064506 (2015)



 $Q_0(\beta)$ is easy to evaluate. Applying the steepest-descent approximation, it corresponds to the situation of a collapsed ring-polymer in the bottom of one of the wells. This results in a simple harmonic vibrational partition function. With $\beta_N = \beta/N$, *N* being the number of beads, one has

$$Q_0(\beta) \simeq \prod_k \frac{1}{\beta_N \hbar \omega_k} = \left(\frac{1}{\beta_N \hbar}\right)^N \frac{1}{\sqrt{\det \mathbf{G}_0}} .$$
 (2)

 ω_k^2 are the eigenvalues of the mass-weighted Hessian of the collapsed ring-polymer, the elements of which are

$$(\mathbf{G}_0)_{ii'} = \frac{2\delta_{ii'} - \delta_{ii'-1} - \delta_{ii'+1}}{(\beta_N \hbar)^2} + \omega_s^2 \delta_{ii'} . \qquad (3)$$

$$Q(\beta) = \sum_{n=0,\text{even}}^{\infty} \frac{2N^n}{n!} Q_n(\beta)$$

Yucheng Zhu, Shuo Yang, Jiaxi Zeng, Wei Fang, Ling Jiang, Donghui Zhang, and Xin-Zheng Li, J. Chem. Phys. 158, 220901 (2023)