



中國人民大學
RENMIN UNIVERSITY OF CHINA

Topological Superradiant Phase in a Degenerate Fermi Gas

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Acknowledgements

Jian-Song Pan (USTC)

Wei Yi (USTC)

Xiong-Jun Liu (PKU)

Dong-Yang Yu (IOP)

Outline

- Introduction
 - Superradiance
 - Spin-orbit coupling
 - Cavity-induced SOC
- Superradiance of a quasi-1D degenerate Fermi gas
 - Superradiance transition
 - Topological properties
 - Phase diagram
 - Experimental signatures
- Summary

Superradiance with two-level emitters

PHYSICAL REVIEW

VOLUME 93, NUMBER 1

JANUARY 1, 1954

Coherence in Spontaneous Radiation Processes

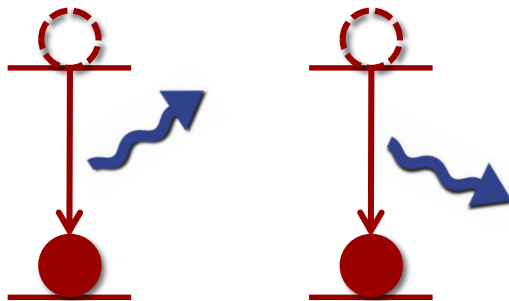
R. H. DICKE

Palmer Physical Laboratory, Princeton University, Princeton, New Jersey

(Received August 25, 1953)

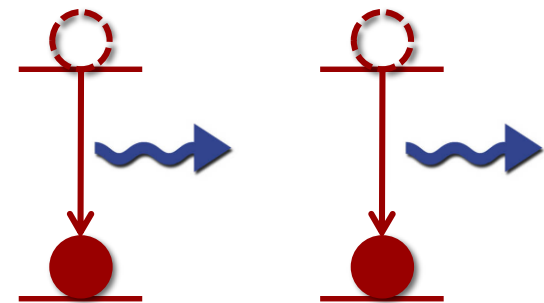
Spontaneous Emission

$$I \sim N$$

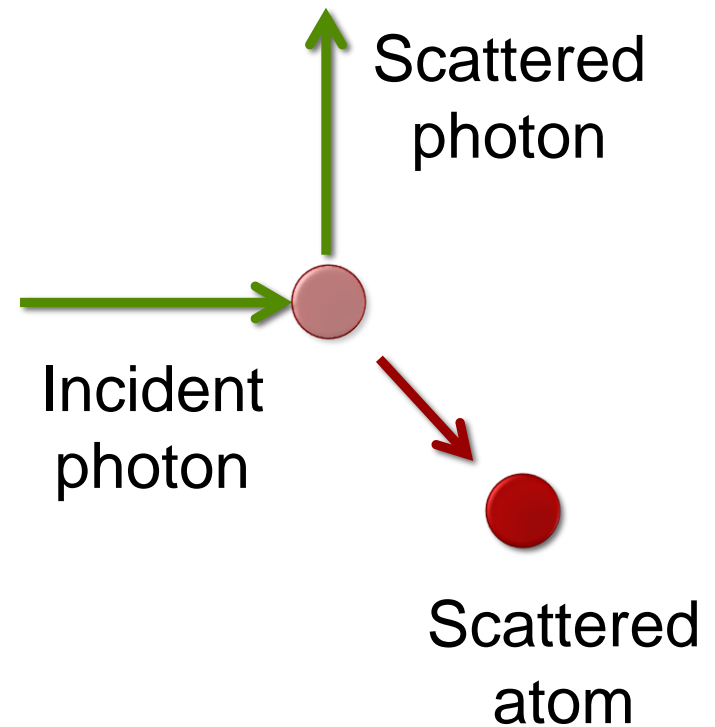
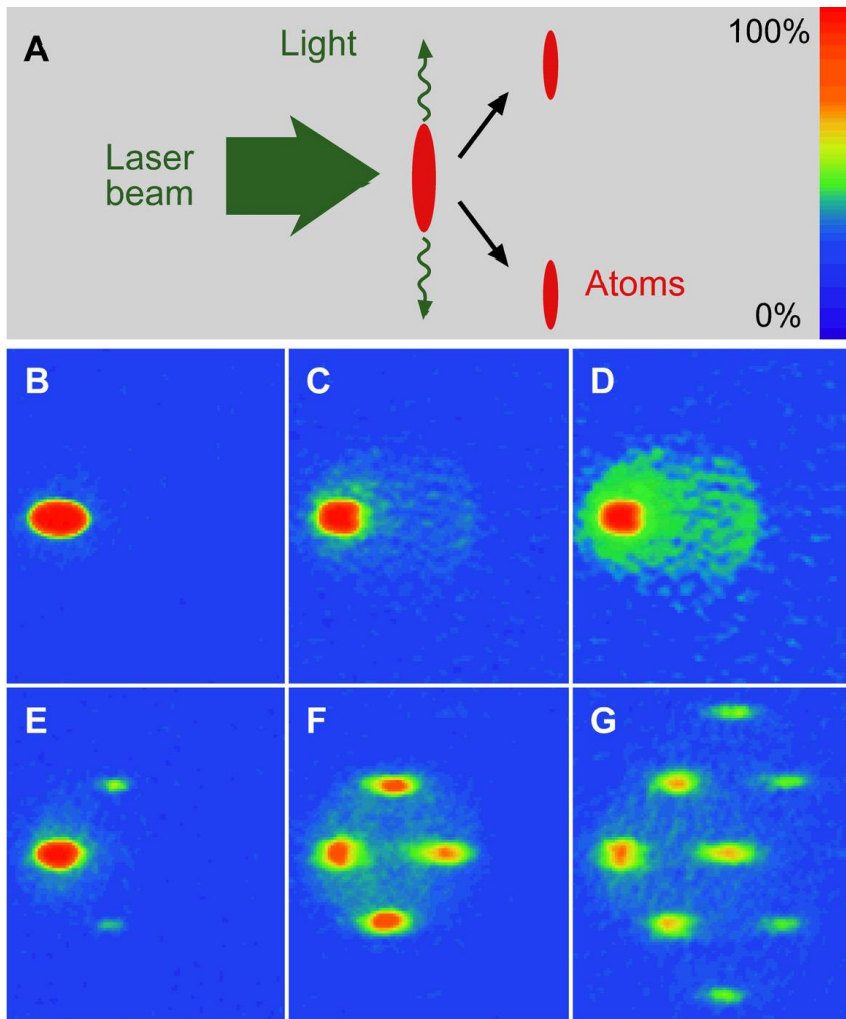


Spontaneous Emission

$$I \sim N^2$$

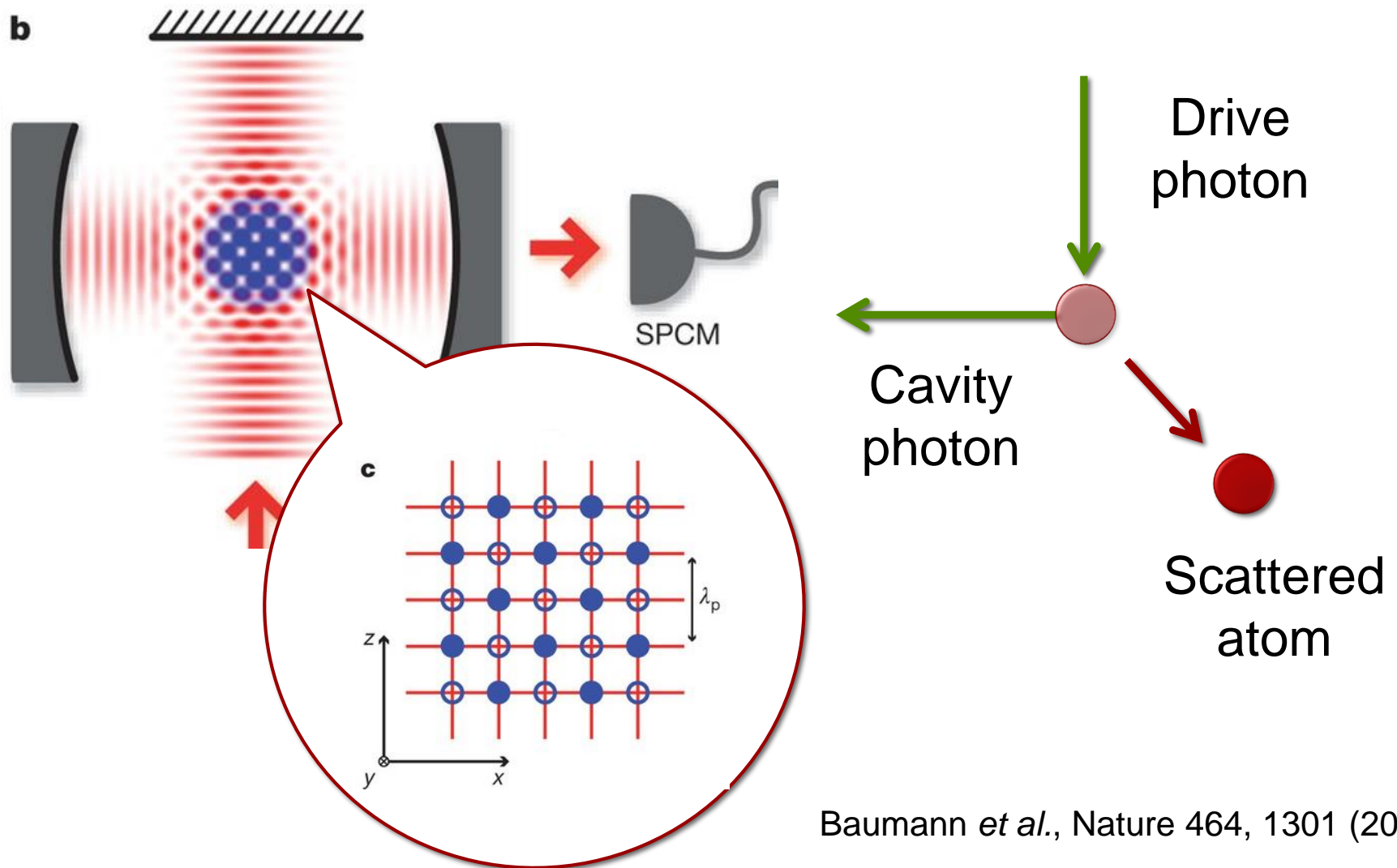


Superradiance with BEC

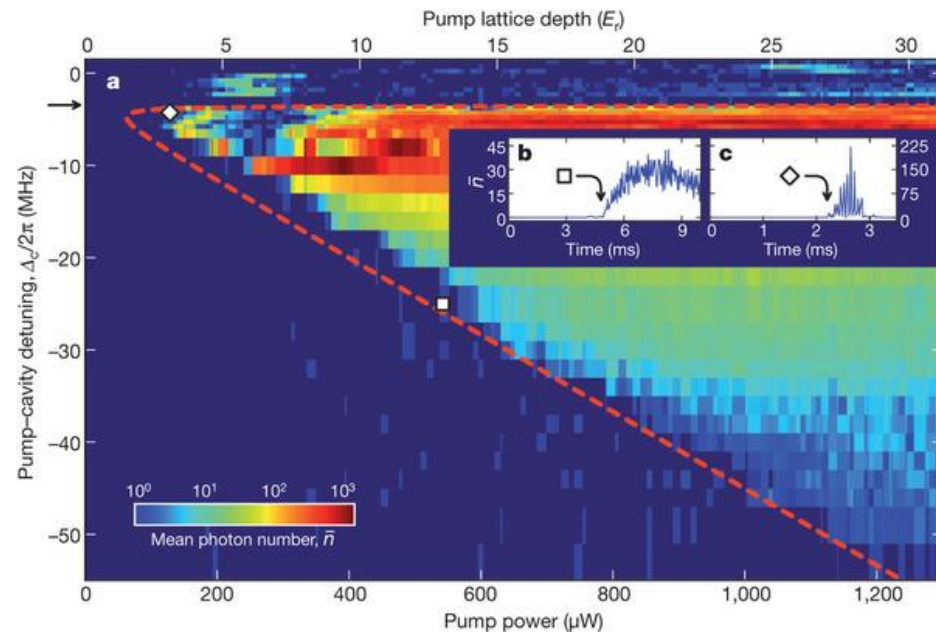
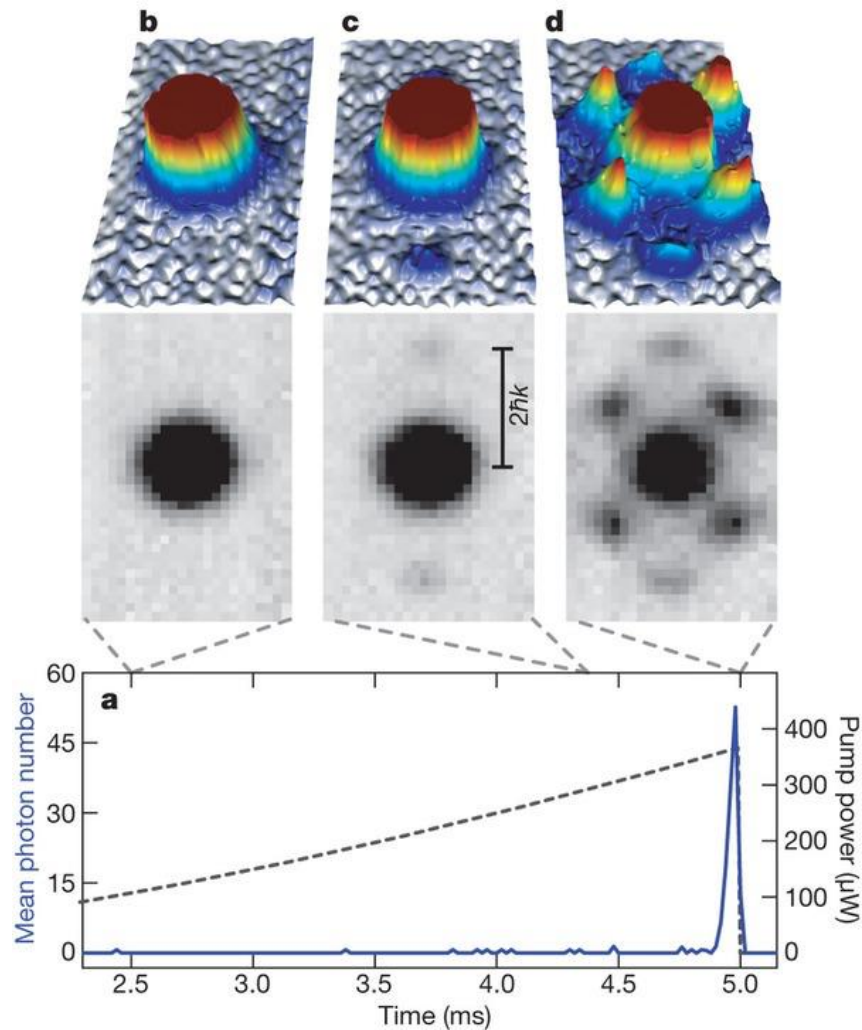


Inouye *et al.*, Science 285, 571 (1999)

Superradiance with BEC in a cavity

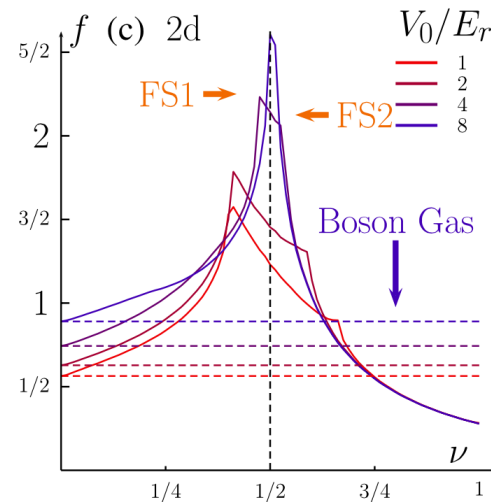
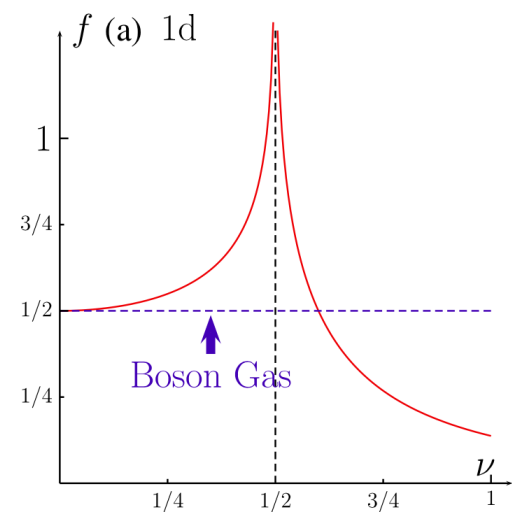
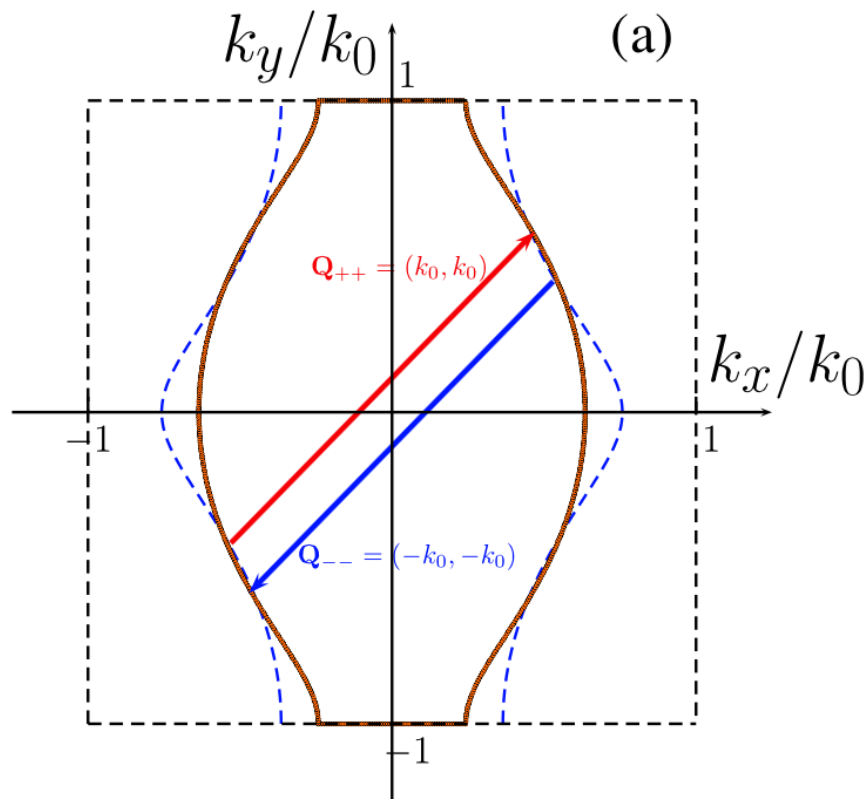


Superradiance with BEC in a cavity



Baumann *et al.*, Nature 464, 1301 (2010)

Superradiance with a degenerate Fermi gas

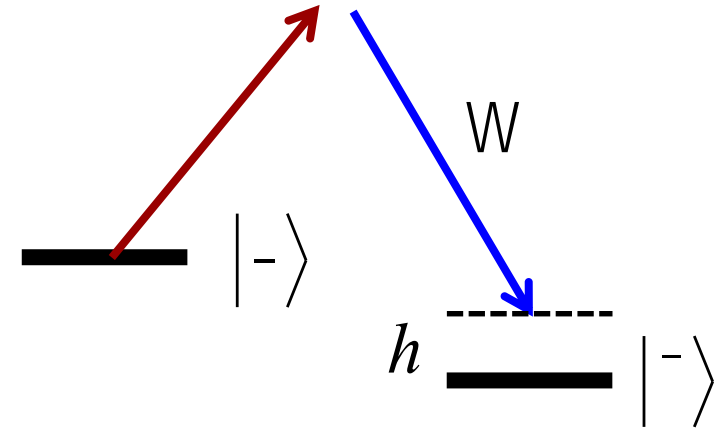


Keeling *et al.*, PRL 112 143002, (2014)
 Piazza *et al.*, PRL 112 143003, (2014)
 Chen *et al.*, PRL 112 143004, (2014)

Synthetic spin-orbit coupling: Raman scheme



Lin et. al., Nature 471, 83 (2011)



$|-\rangle$ $|-\rangle$

$$H = \begin{pmatrix} \frac{k_x^2}{2m} + \frac{h}{2} & \frac{\Omega}{2} e^{i2k_0 x} \\ \frac{\Omega}{2} e^{-i2k_0 x} & \frac{k_x^2}{2m} - \frac{h}{2} \end{pmatrix} \begin{matrix} |-\rangle \\ |-\rangle \end{matrix} \quad U = \begin{pmatrix} e^{-ik_0 x} & 0 \\ 0 & e^{ik_0 x} \end{pmatrix}$$

Synthetic spin-orbit coupling: Raman scheme

$$H_{\text{SO}} = U H U^\dagger = \begin{pmatrix} \frac{(k_x + k_0)^2}{2m} + \frac{\hbar}{2} & \frac{\Omega}{2} \\ \frac{\Omega}{2} & \frac{(k_x - k_0)^2}{2m} - \frac{\hbar}{2} \end{pmatrix}$$

$$= \frac{1}{2m} (k_x + k_0 \sigma_z)^2 + \frac{\Omega}{2} \sigma_x + \frac{\hbar}{2} \sigma_z.$$

- Spin rotation, x to $-z$, z to x

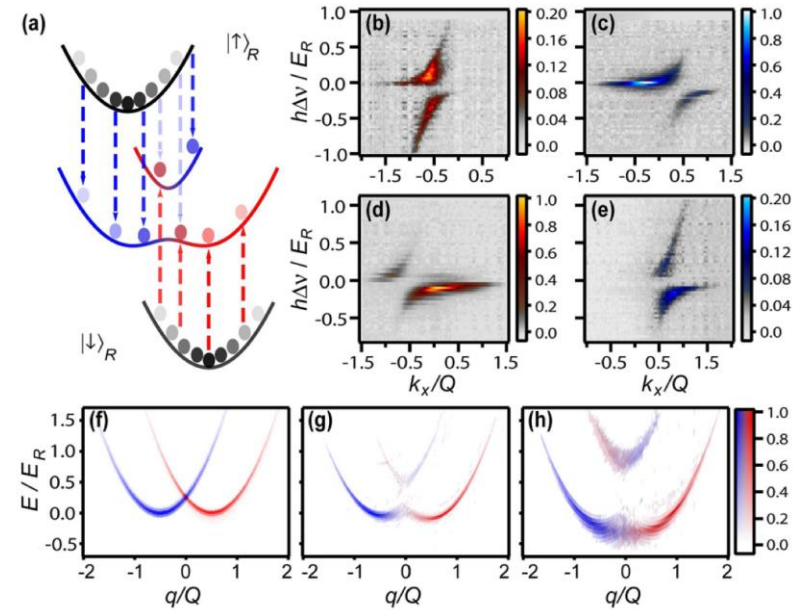
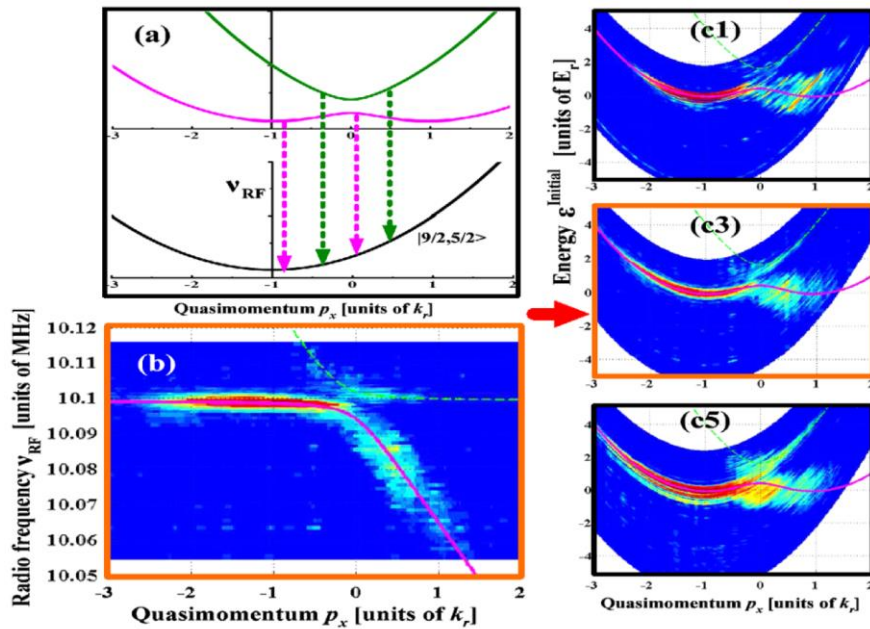
$$H_{\text{SO}} = \frac{1}{2m} \left(k_x^2 + k_0^2 + \boxed{2S_x k_x} \right) - \frac{W}{2} S_z + \frac{\hbar}{2} S_x$$

1D (ERD) SOC

**x-field
2-photon
detuning**

**z-field
Rabi Frequency**

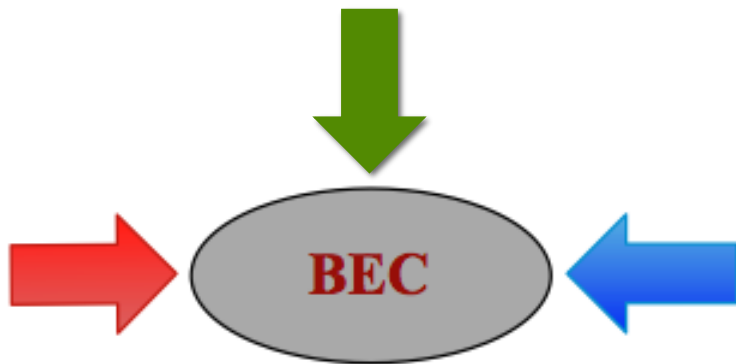
Implementation in Fermi gases



Shanxi University, Taiyuan, China
Wang *et al.*, PRL 109, 095301 (2012)

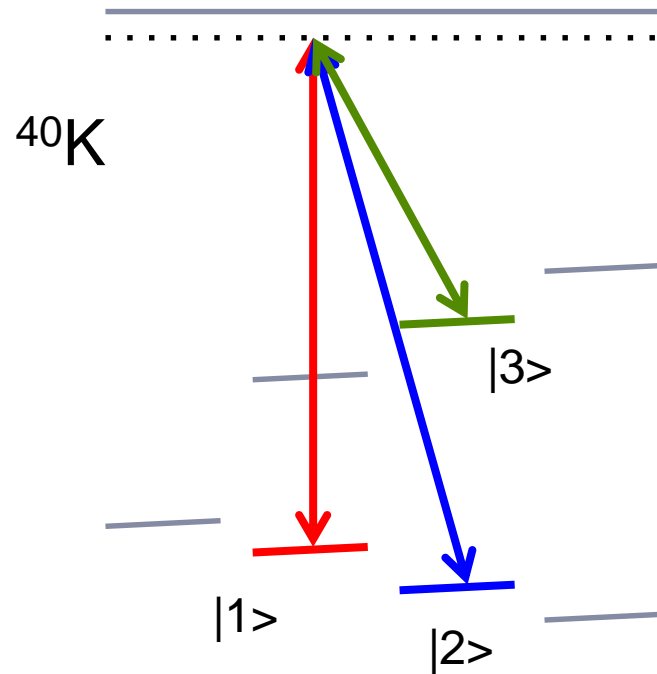
MIT, Boston, USA
Cheuk *et al.*, PRL 109, 095302 (2012)

Implementation of 2D SOC



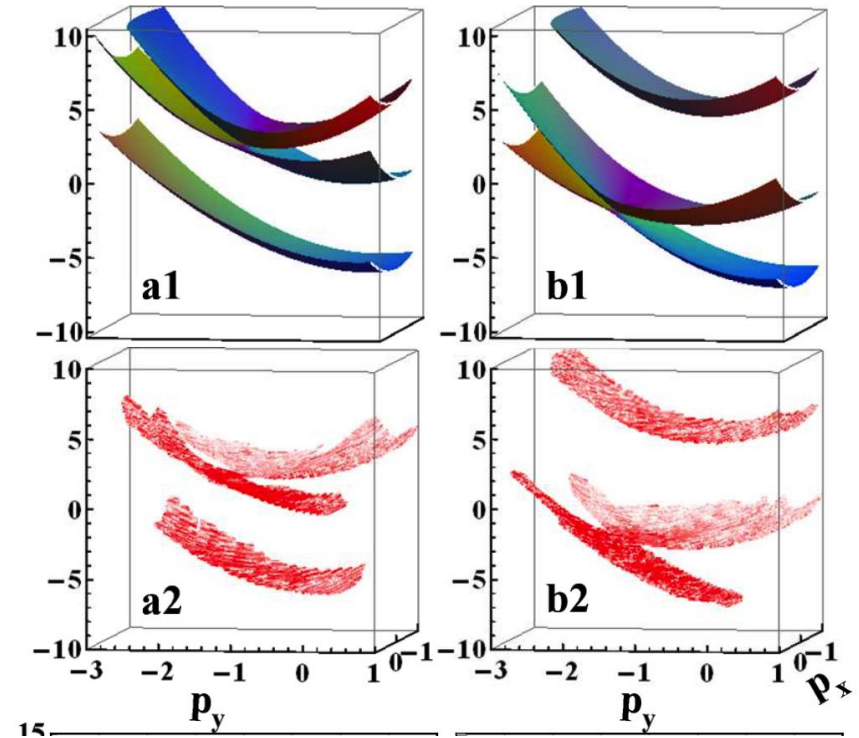
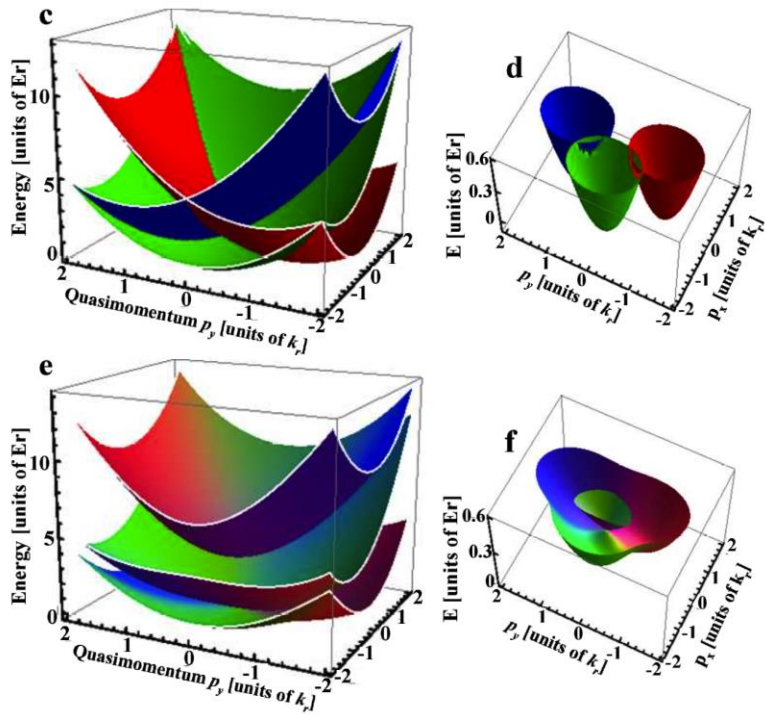
Shanxi University, Taiyuan, China

Huang *et al.*, arXiv:1506.02861 (2015)



$$H = \left(\begin{array}{ccc} (\mathbf{p} - \mathbf{k}_1)^2 + d_1 & -W_{12} / 2 & -W_{13} / 2 \\ -W_{12} / 2 & (\mathbf{p} - \mathbf{k}_2)^2 + d_2 & -W_{23} / 2 \\ -W_{13} / 2 & -W_{23} / 2 & (\mathbf{p} - \mathbf{k}_3)^2 + d_3 \end{array} \right) + k_z^2$$

Implementation of 2D SOC



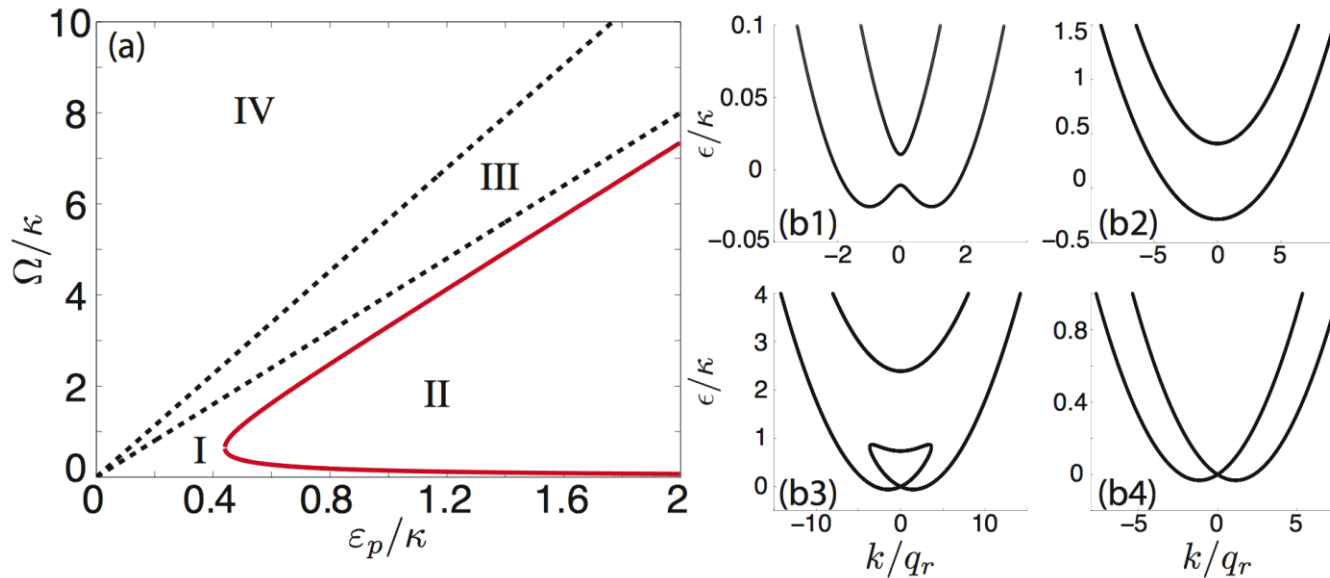
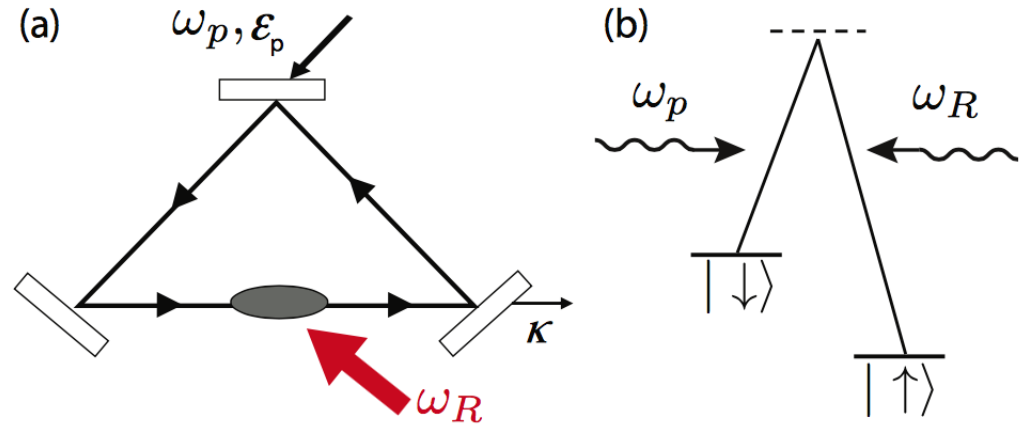
Shanxi University, Taiyuan, China
Huang *et al.*, arXiv:1506.02861 (2015)

$$H_{\text{SO}} = a_x s_x k_x + a_y s_y k_y + a_z s_z k_z$$

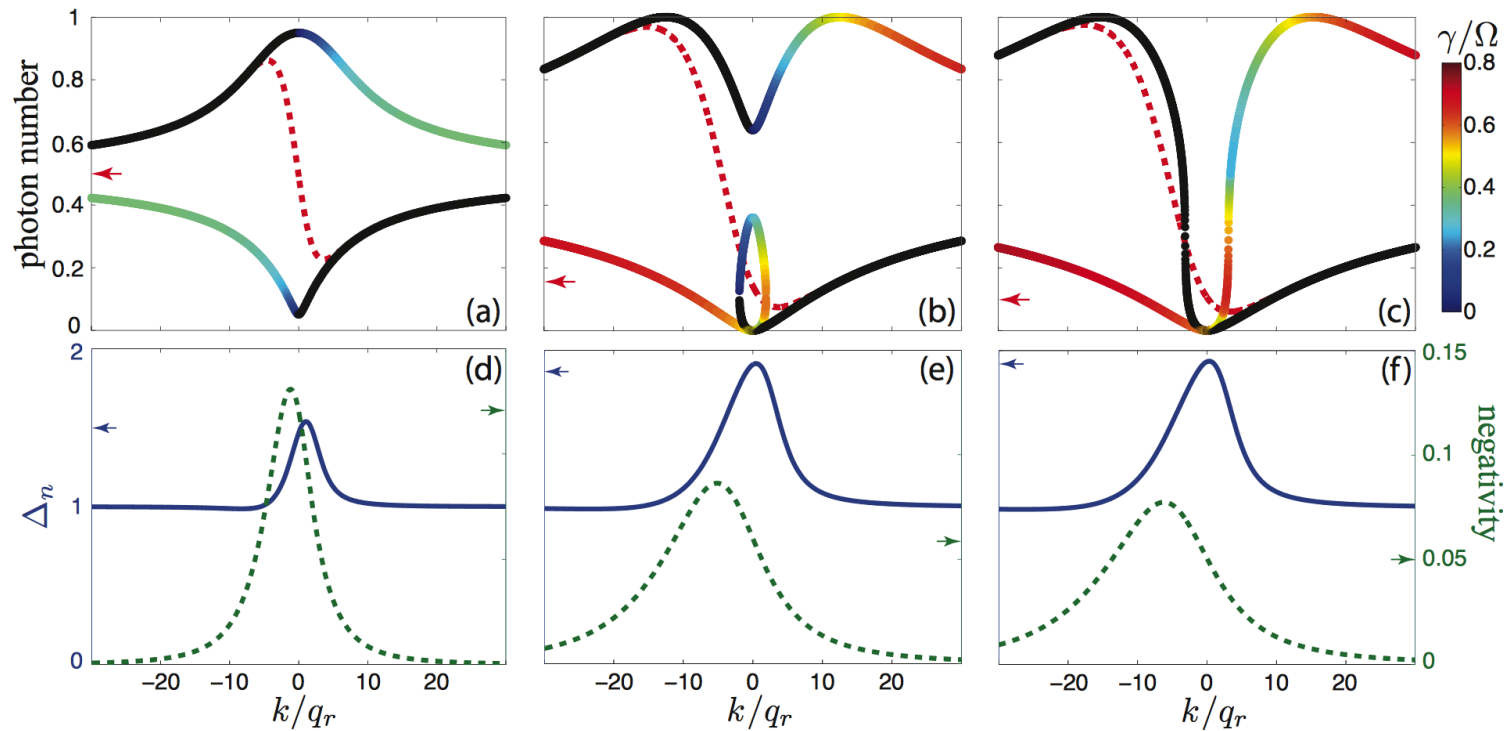
$a_x = a_y = a_z$	3D isotropic SOC	$a\vec{s} \times \vec{k}$
$a_x = a_y$ $a_z = 0$	Rashba SOC	$a(s_x k_x + s_y k_y)$
$a_x = -a_y$ $a_z = 0$	Dresselhaus SOC	$a(s_x k_x - s_y k_y)$
$a_y = a_z = 0$	1D SOC Equal R-D SOC	$a s_x k_x$

Cavity-induced SOC

Dong, Zhou, Wu, Ramachandhran, Pu
 PRA 89, 011602(R) (2014);
 Dong, Zhu, Pu, Atom 3, 183 (2015)

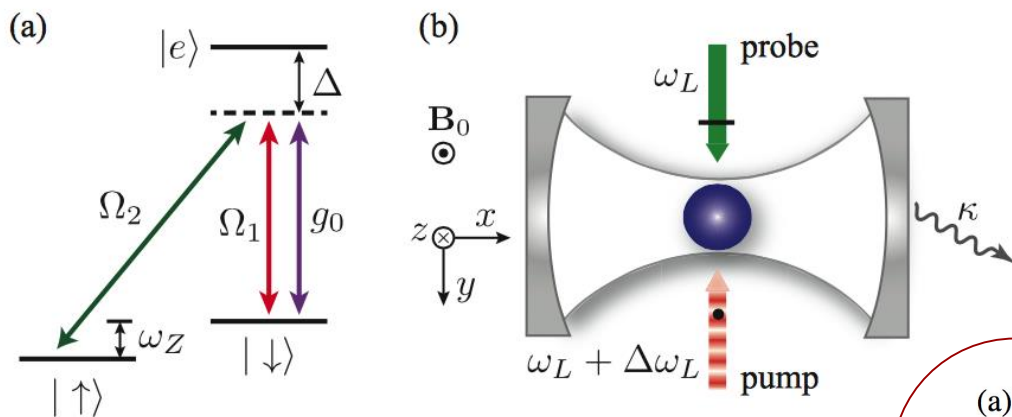


Cavity-induced SOC

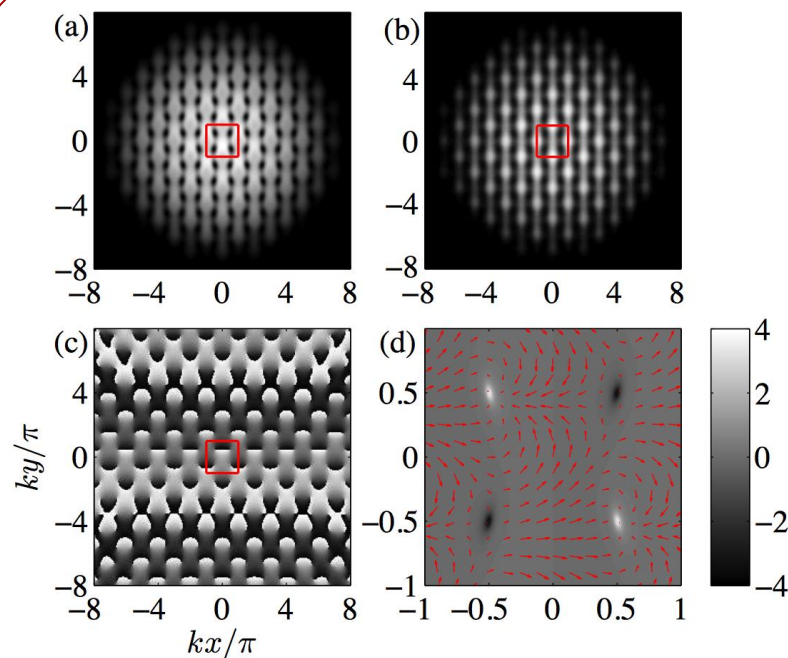
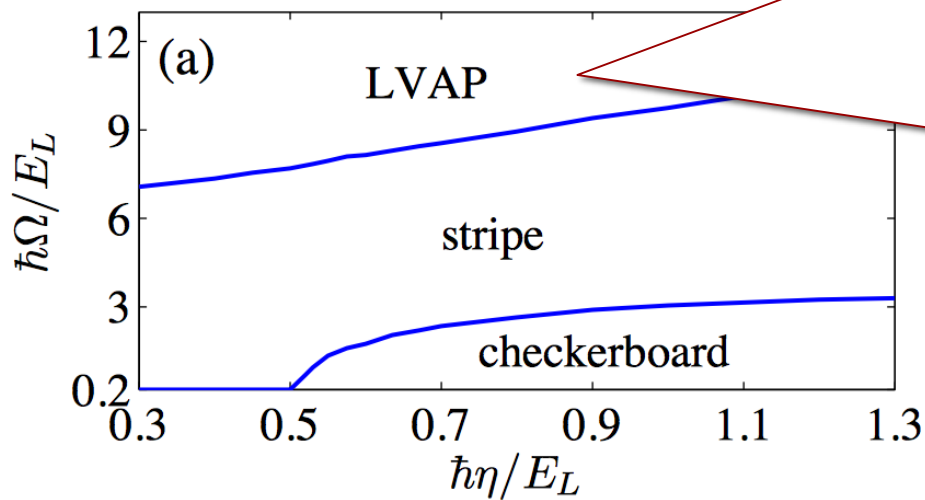


Dong, Zhu, Pu, Atom 3, 183 (2015)

Cavity-induced SOC



Deng, Cheng, Jing, Yi,
PRL 112, 143007 (2014)

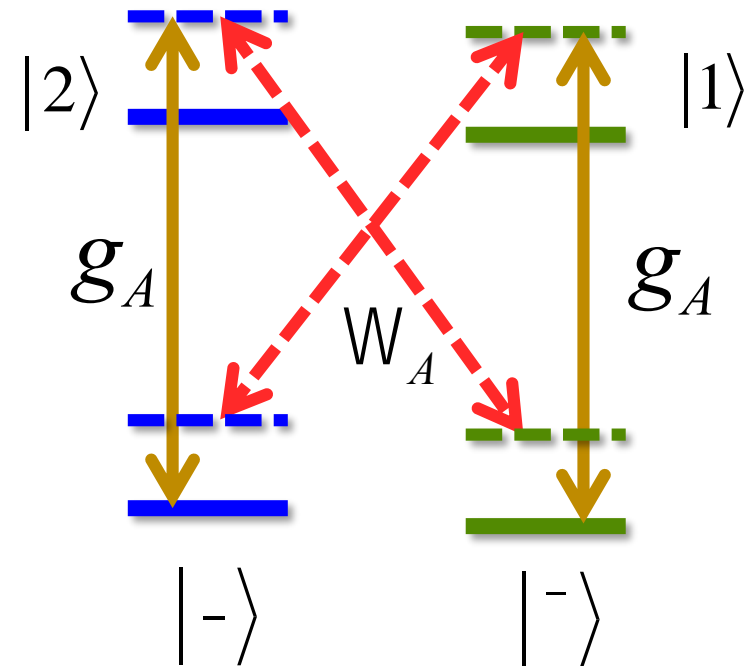
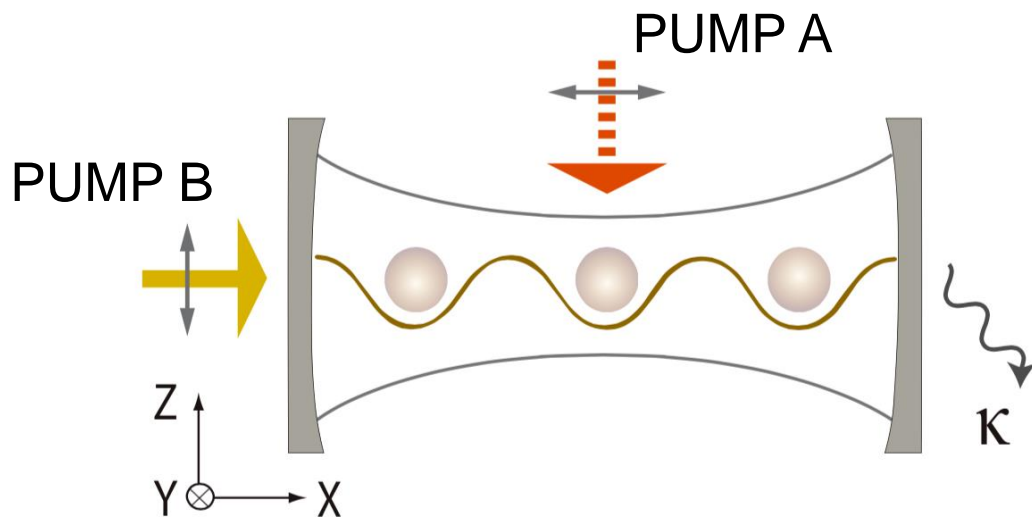


Outline

- Introduction
 - Superradiance of degenerate quantum gases
 - Spin-orbit coupling
 - Cavity-induced SOC
- Superradiance of a 1D degenerate Fermi gas
 - Superradiance transition
 - Topological properties
 - Phase diagram
 - Experimental signatures
- Summary

System

- A quasi-1D two-component Fermi gas in a cavity

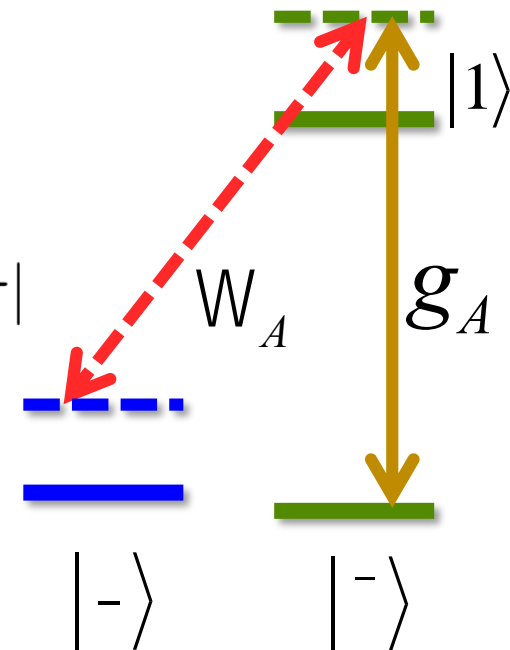


Effective Hamiltonian

- Dynamics of a single particle $\mathcal{H} = \mathcal{H}_0 + \mathcal{H}_I$

$$\mathcal{H}_0 = \hbar\omega_c \hat{a}^\dagger \hat{a} + \sum_{j=1,2} \hbar\omega_j |j\rangle\langle j| + \sum_{\sigma=\uparrow,\downarrow} \hbar\omega_\sigma |\sigma\rangle\langle\sigma|$$

$$\mathcal{H}_I = - \left(\Omega_{1\uparrow} \sigma_{1\uparrow}^+ + \Omega_{2\downarrow} \sigma_{2\downarrow}^+ \right) e^{-i(\omega_A t + k_0 z)} \\ - \left(g_{1\downarrow} \sigma_{1\downarrow}^+ \hat{a} + g_{2\uparrow} \sigma_{2\uparrow}^+ \hat{a} \right) \cos(k_0 x) + h.c.$$



- Unitary transformation $U(t) = \exp \left[i \left(\sum_{j=1,2} |j\rangle\langle j| + \hat{a}^\dagger \hat{a} \right) \omega_A t \right]$

$$\mathcal{H}_{eff} = - \left[\Delta_A - \xi_A \cos^2(k_0 x) \right] \hat{a}^\dagger \hat{a} + \hbar m_z \sigma_z \\ + \eta \left[(\hat{a} e^{ik_0 z} + \hat{a}^\dagger e^{-ik_0 z}) | \uparrow \rangle \langle \downarrow | + H.C. \right] \cos(k_0 x)$$

Effective Hamiltonian

- Transversal degrees of freedom

$$\hat{H} = \sum_{\sigma} \int d\mathbf{r} \hat{\Psi}_{\sigma}^{\dagger} \left[\frac{\mathbf{p}^2}{2m} + U(\mathbf{r}) + (V_0 + \xi_A \hat{a}^{\dagger} \hat{a}) \cos^2(k_0 x) + \xi_{\sigma} m_z \right] \hat{\Psi}_{\sigma} - \Delta_A \hat{a}^{\dagger} \hat{a} \\ + \eta \left[\int d\mathbf{r} \hat{\Psi}_{\uparrow}^{\dagger} (\hat{a} e^{ik_0 z} + \hat{a}^{\dagger} e^{-ik_0 z}) \cos(k_0 x) \hat{\Psi}_{\downarrow} + h.c. \right].$$

- 1D model

$$\hat{H} = \sum_{\sigma} \int dx \hat{\psi}_{\sigma}^{\dagger} \left[\frac{p_x^2}{2m} + (V_0 + \xi_A \hat{a}^{\dagger} \hat{a}) \cos^2(k_0 x) + \xi_{\sigma} m_z \right] \hat{\psi}_{\sigma} \\ + \eta_A (\hat{a} + \hat{a}^{\dagger}) \left[\int dx \hat{\psi}_{\uparrow}^{\dagger} \cos(k_0 x) \hat{\psi}_{\downarrow} + h.c. \right],$$

Effective Hamiltonian

- Equation of motion

$$i\dot{\hat{a}} = \hat{a} \left[\xi_A \sum_{\sigma} \int dx \hat{\psi}_{\sigma}^{\dagger} \cos^2(k_0 x) \hat{\psi}_{\sigma} - \Delta_A - i\kappa \right] \\ + \eta_A \left[\int dx \hat{\psi}_{\downarrow}^{\dagger} \cos(k_0 x) \hat{\psi}_{\uparrow} + h.c. \right]$$

- Steady state solution

$$\alpha \approx \frac{\eta_A \int dx \cos(k_0 x) \left[\langle \hat{\psi}_{\downarrow}^{\dagger} \hat{\psi}_{\uparrow} \rangle + h.c. \right]}{\Delta_A + i\kappa - \xi_A \sum_{\sigma} \int dx \langle \hat{\psi}_{\sigma}^{\dagger} \hat{\psi}_{\sigma} \rangle \cos^2(k_0 x)}$$

Effective Hamiltonian

OL generated by
Longitudinal
Pumping

Cavity Field

Effective Zeeman
Splitting

$$\hat{H} = \sum_{\sigma} \int dx \hat{\psi}_{\sigma}^{\dagger} \left[\frac{p_x^2}{2m} + \boxed{V_0} + \xi_A \boxed{\hat{a}^{\dagger} \hat{a}} \cos^2(k_0 x) + \xi_{\sigma} \boxed{m_z} \right] \hat{\psi}_{\sigma}$$

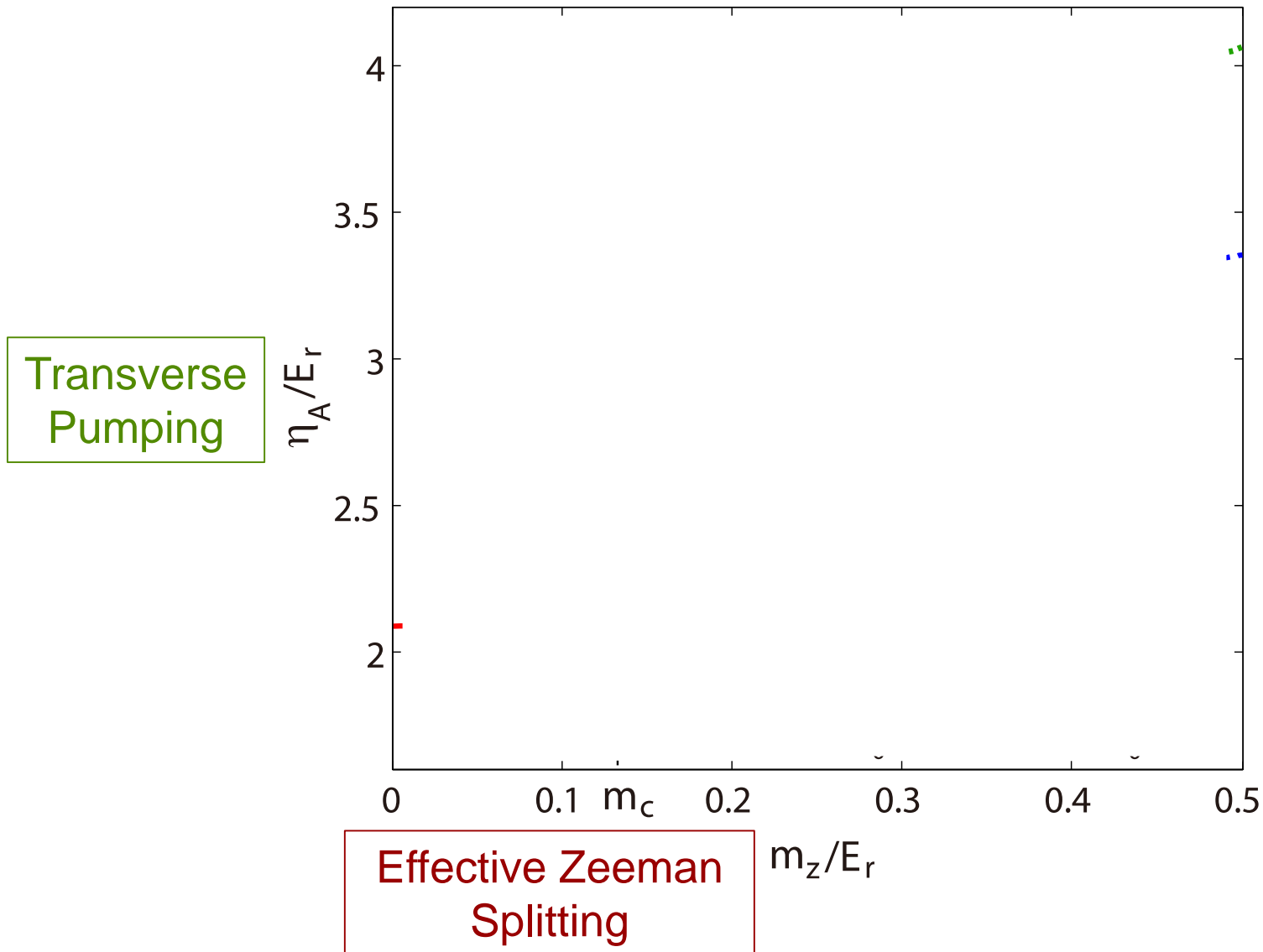
$$+ \boxed{\eta_A} (\alpha^* + \alpha) \left[\int dx \hat{\psi}_{\uparrow}^{\dagger} \cos(k_0 x) \hat{\psi}_{\downarrow} + \text{H.C.} \right]$$

Transverse Pumping

- Solve with

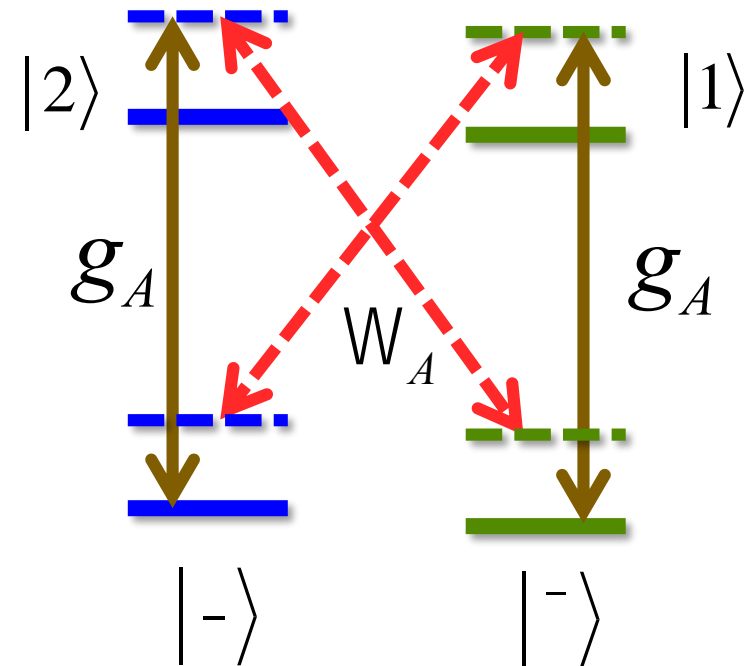
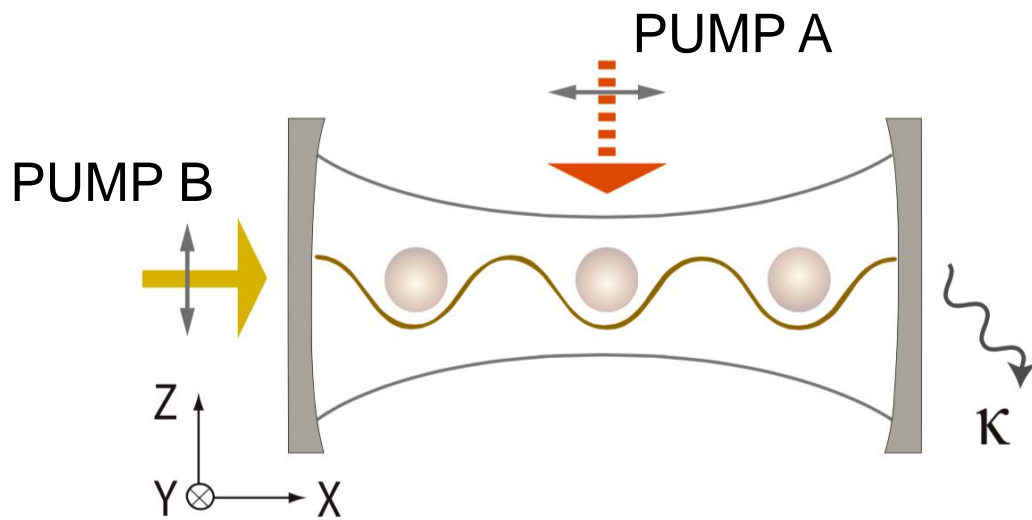
$$\alpha \approx \frac{\eta_A \int dx \cos(k_0 x) \left[\langle \hat{\psi}_{\downarrow}^{\dagger} \hat{\psi}_{\uparrow} \rangle + h.c. \right]}{\Delta_A + i\kappa - \xi_A \sum_{\sigma} \int dx \langle \hat{\psi}_{\sigma}^{\dagger} \hat{\psi}_{\sigma} \rangle \cos^2(k_0 x)}$$

Phase diagram



System

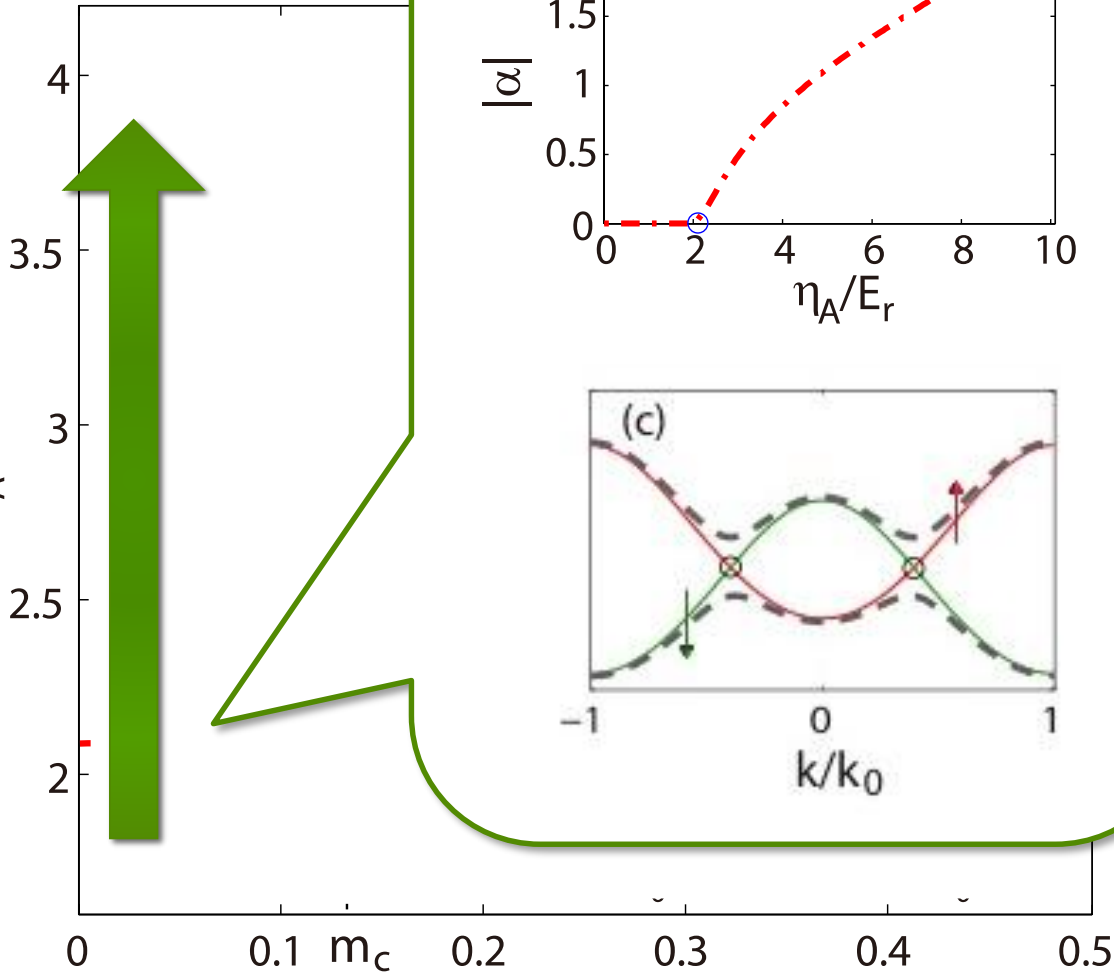
- A quasi-1D two-component Fermi gas in a cavity



Phase diagram

Transverse Pumping

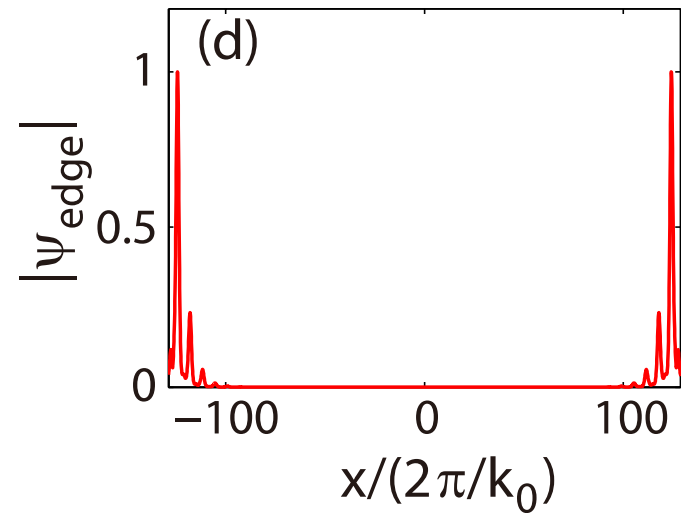
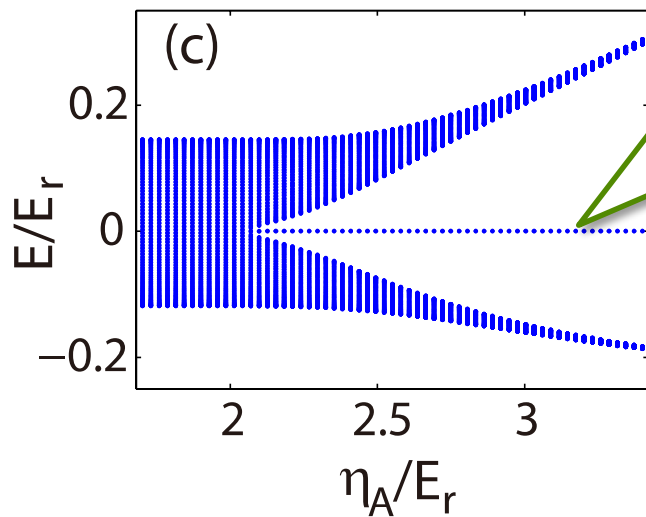
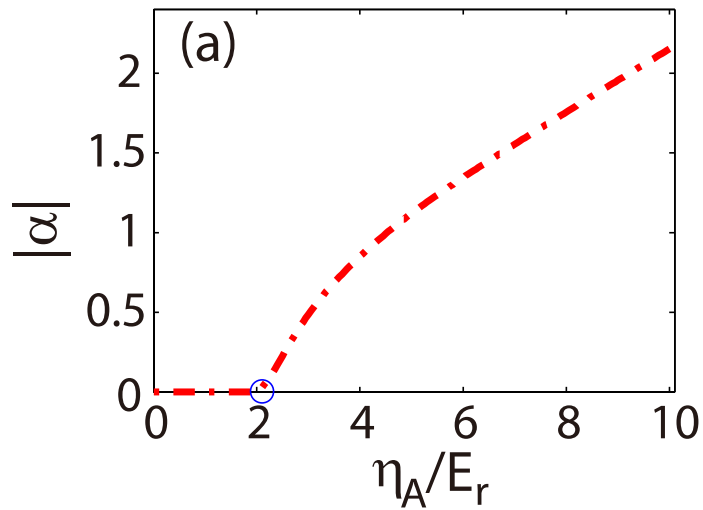
η_A/E_r



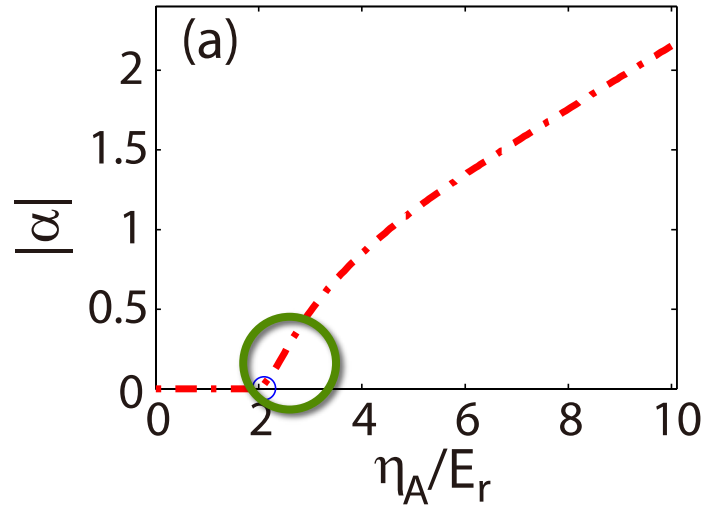
Effective Zeeman Splitting

m_z/E_r

Topological properties



Topological properties



$$\hat{H}_{TI} = - \sum_{\langle i,j \rangle, \sigma} t_s^{1,1}(i,j) \hat{\psi}_{i\sigma}^\dagger \hat{\psi}_{j\sigma} + \sum_{\langle i,j \rangle} \left(t_{so}^{1,1}(i,j) \hat{\psi}_{i\uparrow}^\dagger \hat{\psi}_{j\downarrow} + \text{H.C.} \right) + m_z \sum_{i,\sigma} \xi_\sigma \hat{\psi}_{i\sigma}^\dagger \hat{\psi}_{i\sigma},$$

$$t_s^{n,n'}(i,j) = \int dx \phi_{ni}^*(x) \left[\frac{\hat{p}_x^2}{2m} + \left(V_0 + \xi_A |\alpha|^2 \right) \cos^2(k_0 x) \right] \phi_{n'j}(x),$$

$$t_{so}^{n,n'}(i,j) = \eta_A (\alpha^* + \alpha) \int dx \phi_{ni}^*(x) \cos(k_0 x) \phi_{n'j}(x),$$

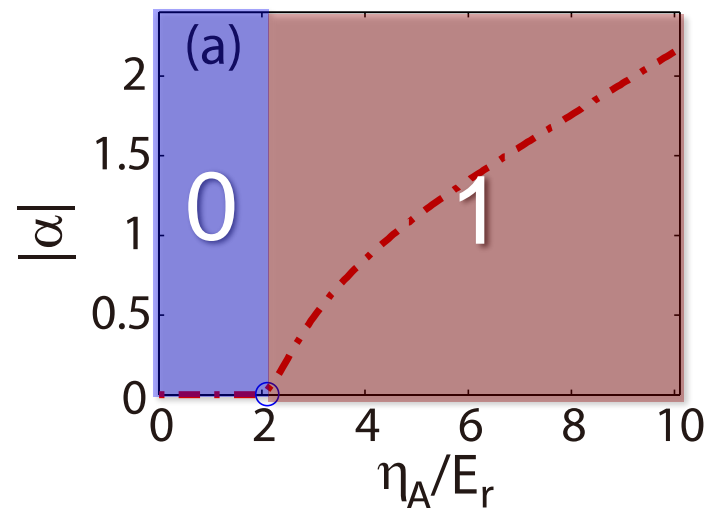
Topological properties

$$\hat{H}_{TI} = \sum_{k \in BZ} \begin{pmatrix} \hat{\psi}_{k\uparrow}^\dagger & \hat{\psi}_{k\downarrow}^\dagger \end{pmatrix} [h_y(k) \sigma_y + h_z(k) \sigma_z] \begin{pmatrix} \hat{\psi}_{k\uparrow} \\ \hat{\psi}_{k\downarrow} \end{pmatrix},$$

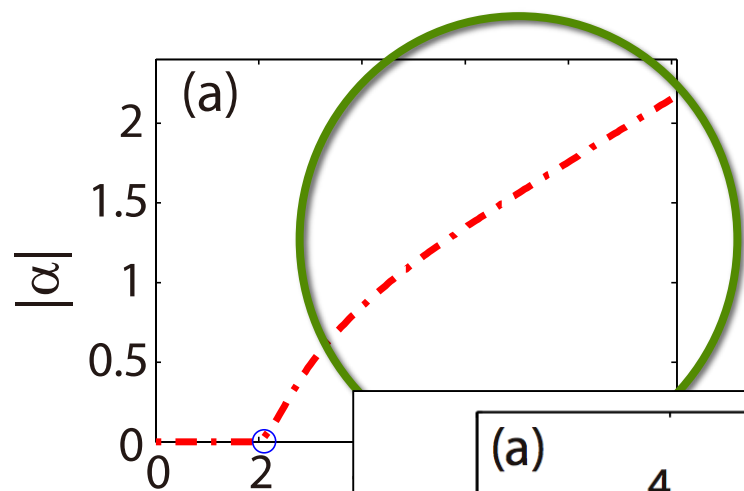
$$h_y(k) = 2t_{so} \sin(ka)$$

$$h_z(k) = m_z - 2t_s \cos(ka)$$

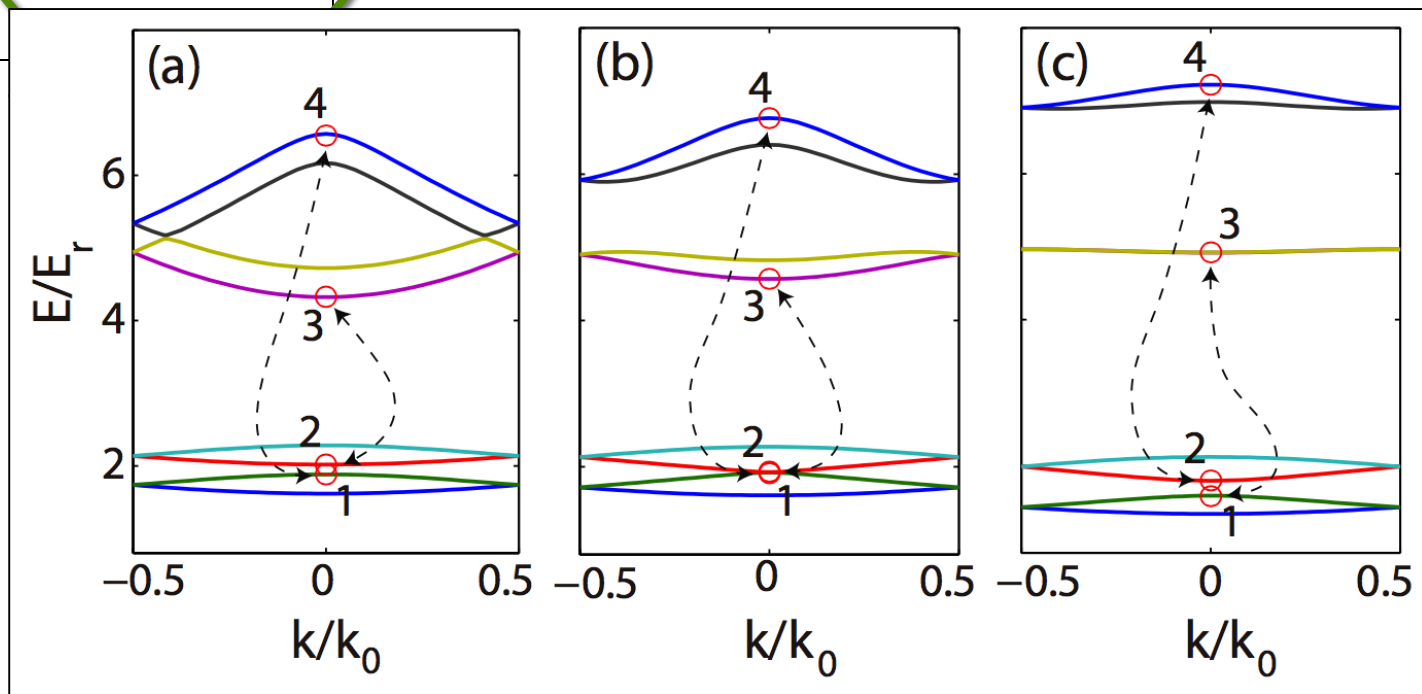
$$\mathcal{W} = \sum_{\nu, \nu' = y, z} \oint \frac{dk}{4\pi} \epsilon_{\nu, \nu'} \hat{h}_\nu^{-1}(k) \partial_k \hat{h}_{\nu'}(k),$$



Topological properties

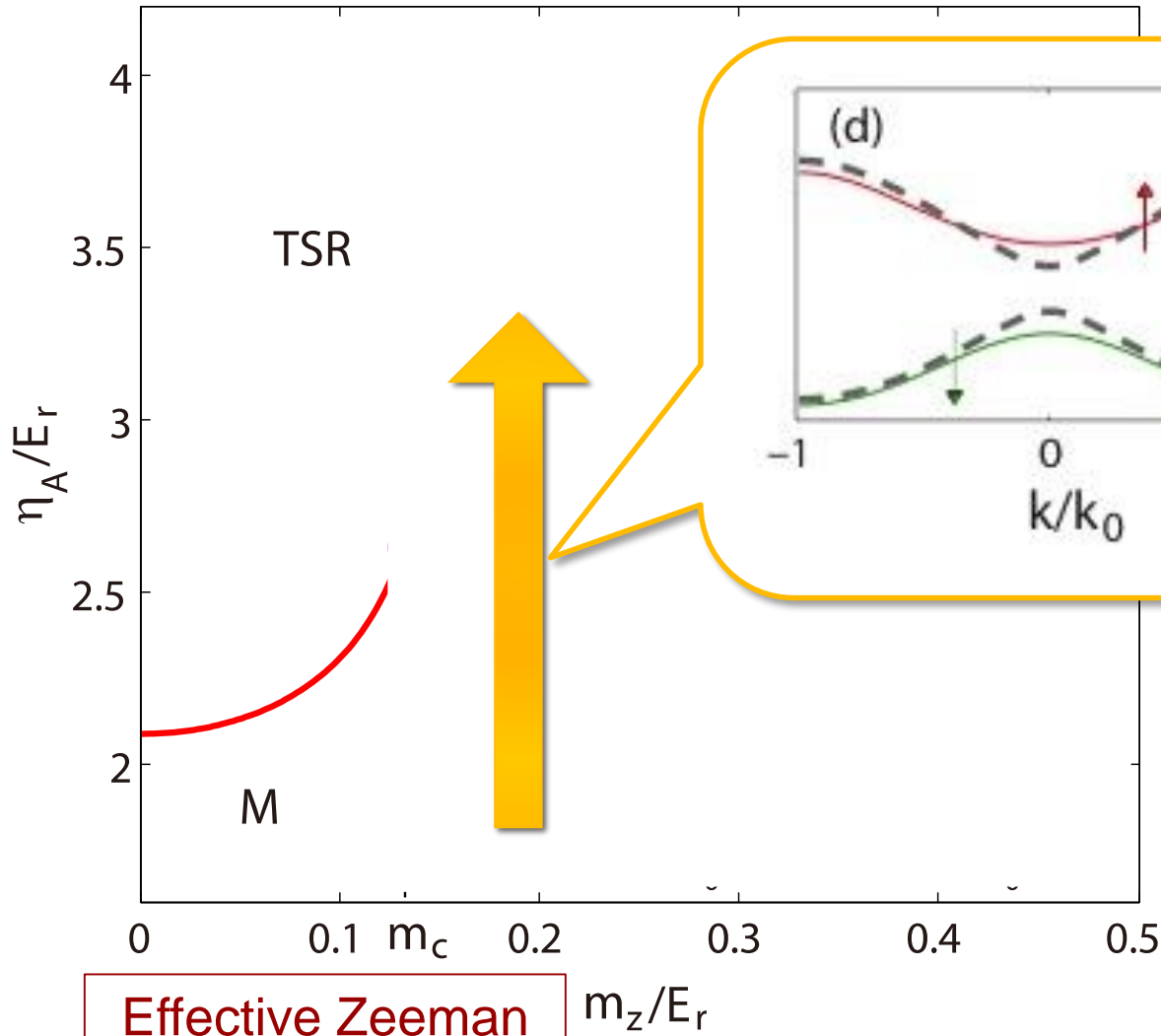


See also, Zhou, Cui., arXiv:1507.01341



Phase diagram

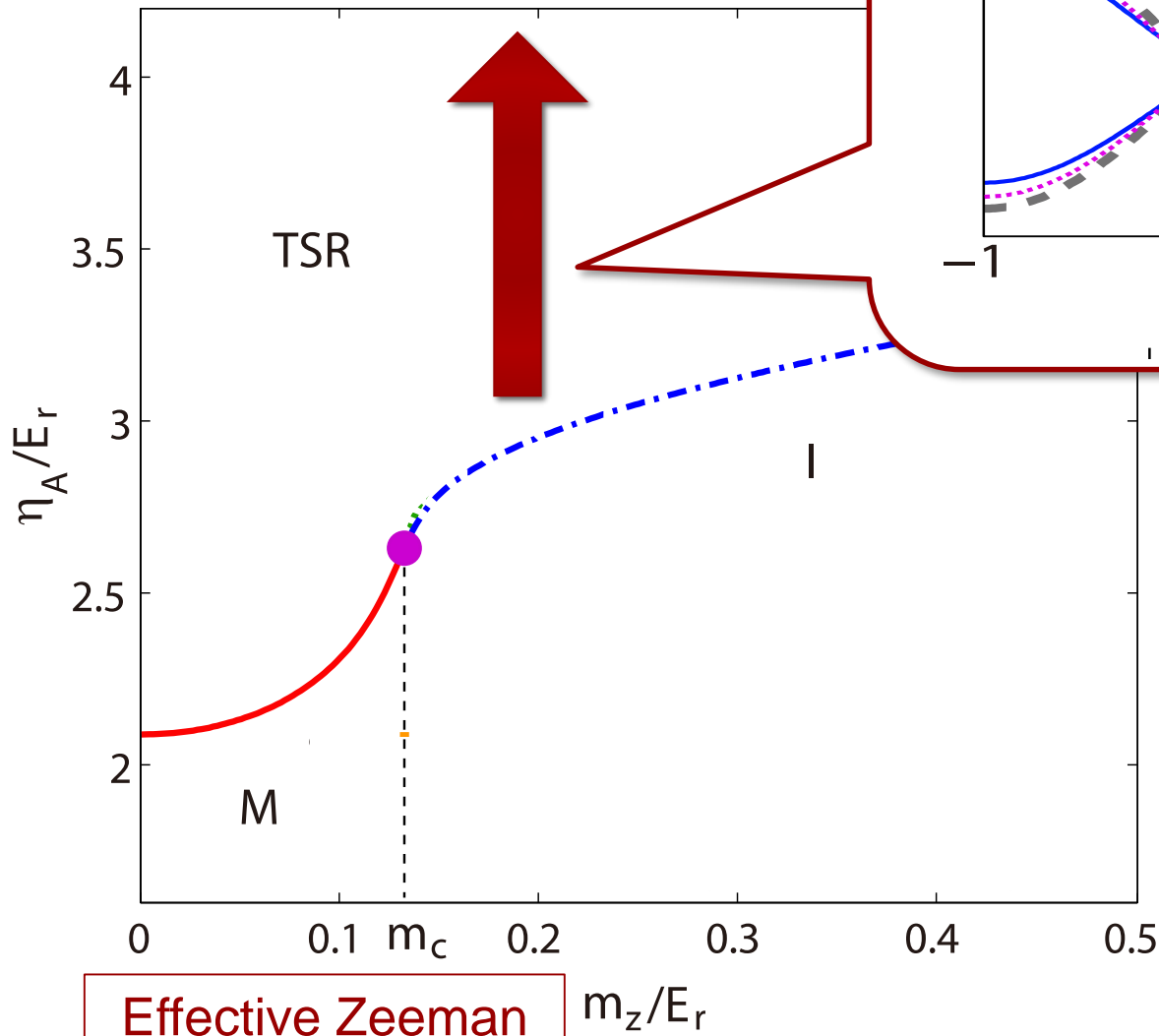
Transverse Pumping



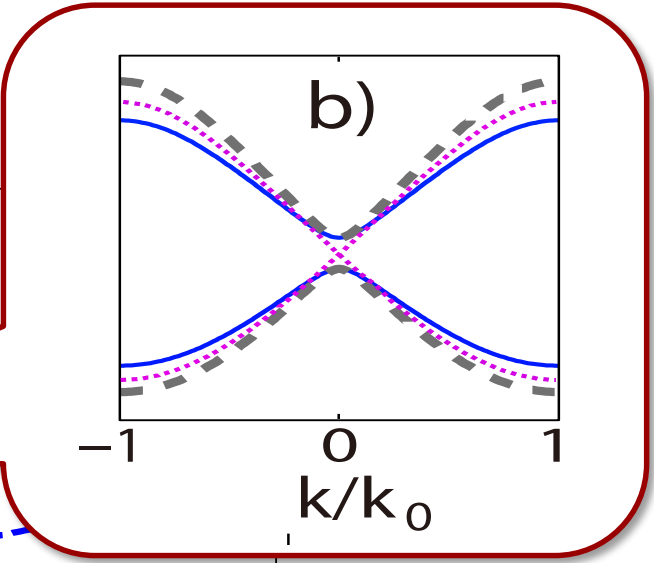
Effective Zeeman Splitting

Phase diagram

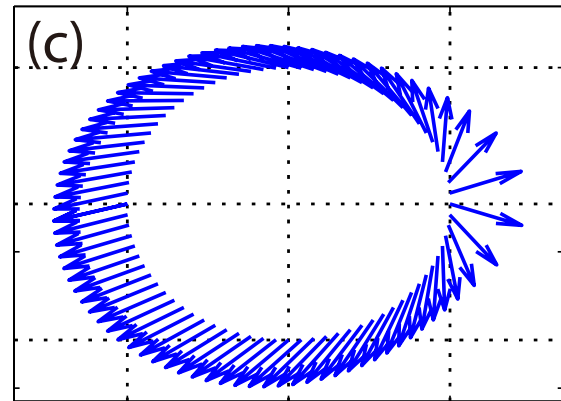
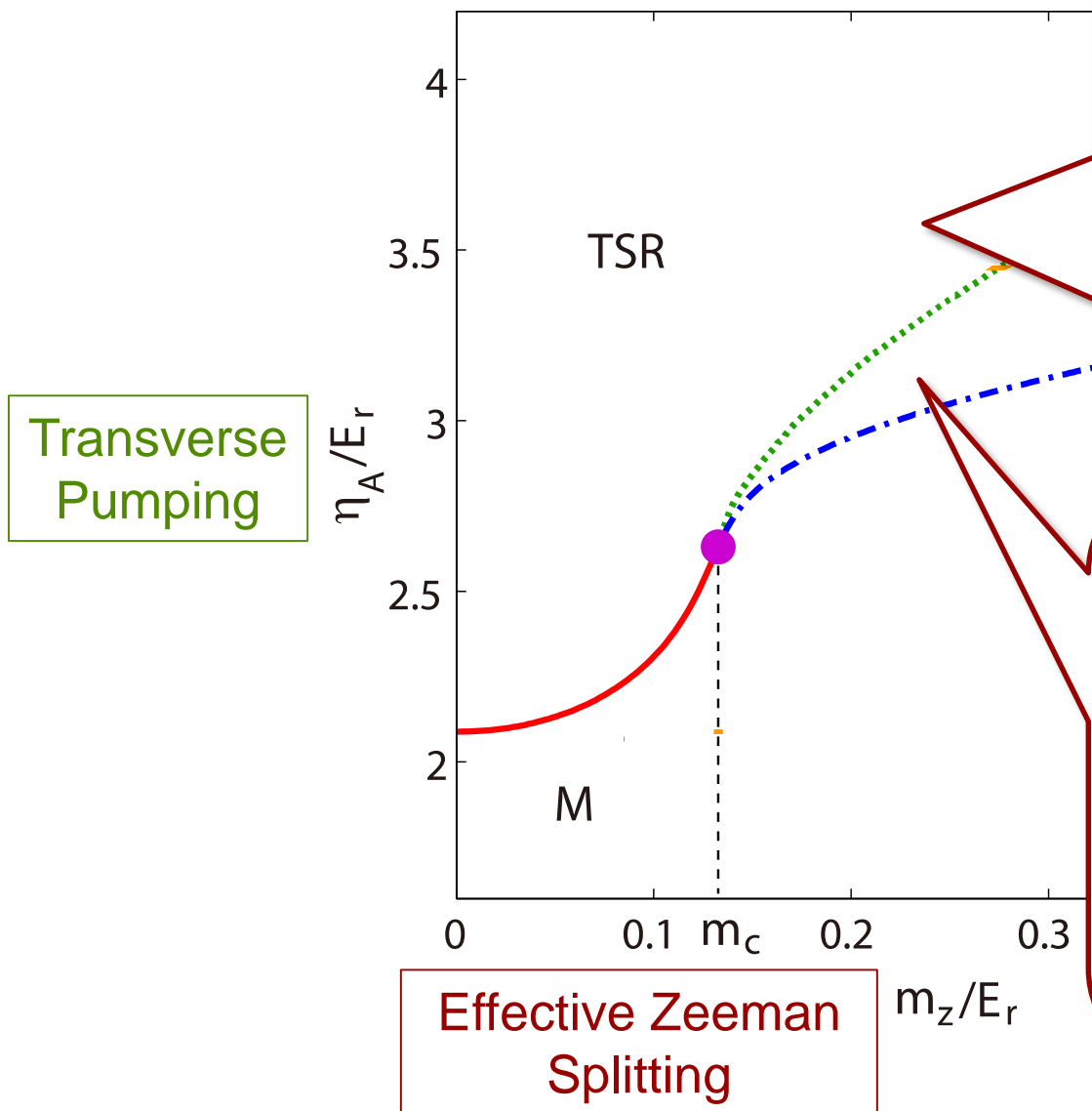
Transverse Pumping



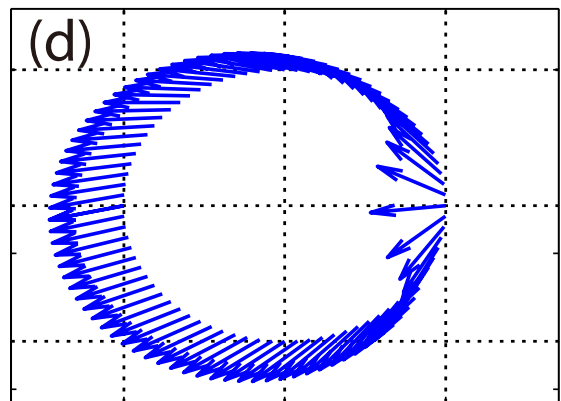
Effective Zeeman Splitting



Phase diagram



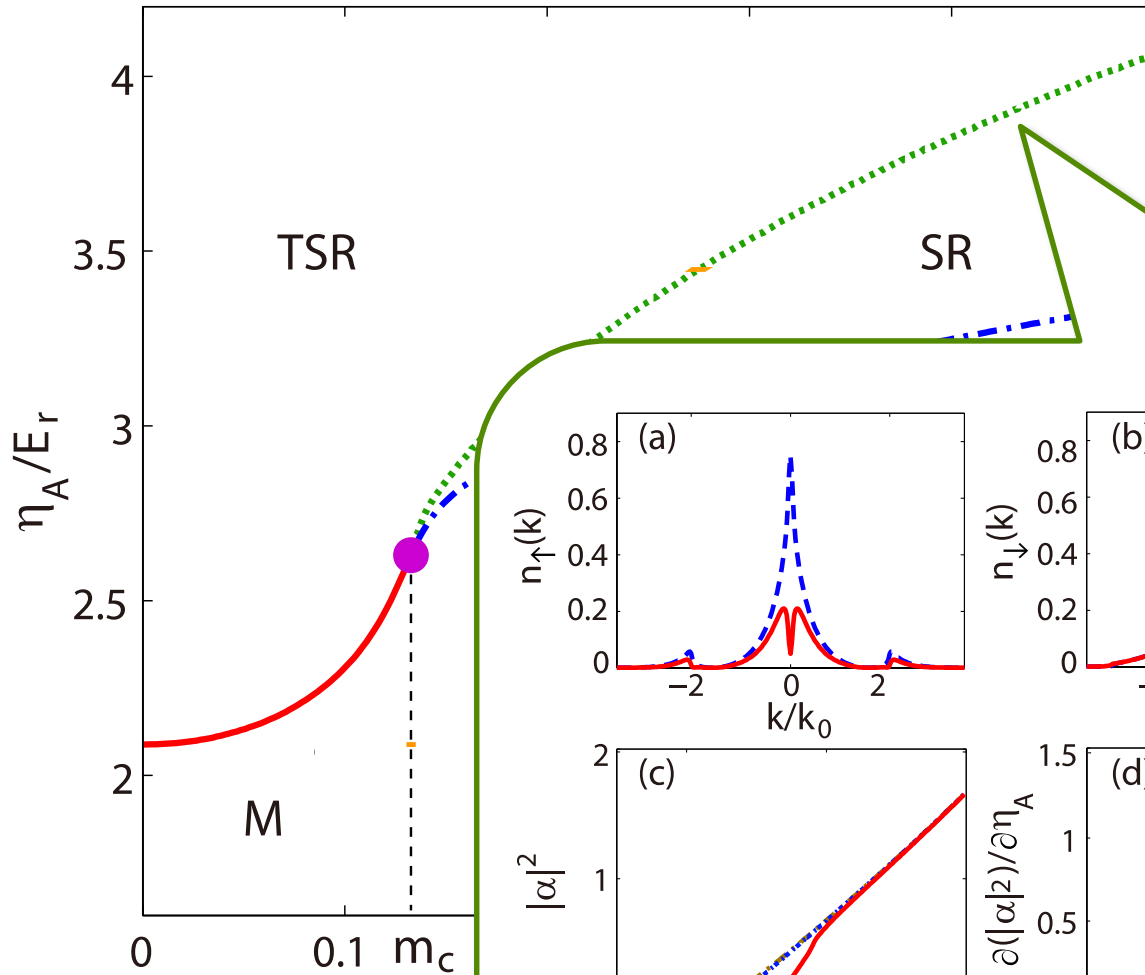
$$\langle S_z(k) \rangle$$



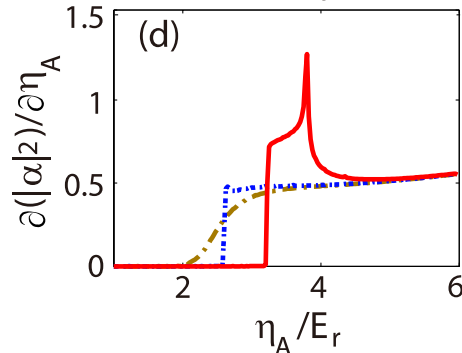
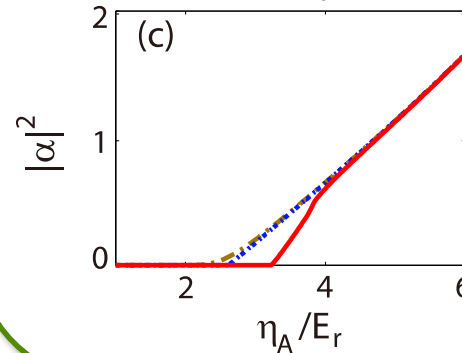
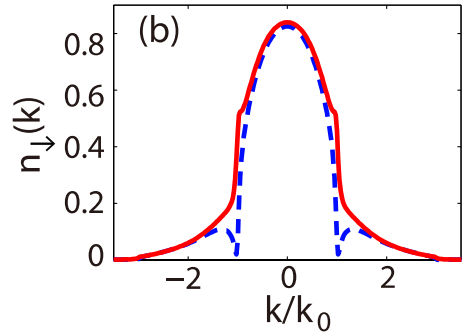
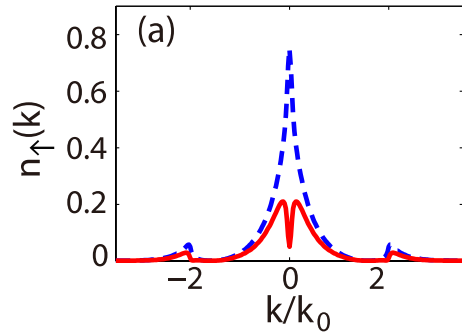
$$\langle S_z(k) \rangle$$

Observables

Transverse Pumping

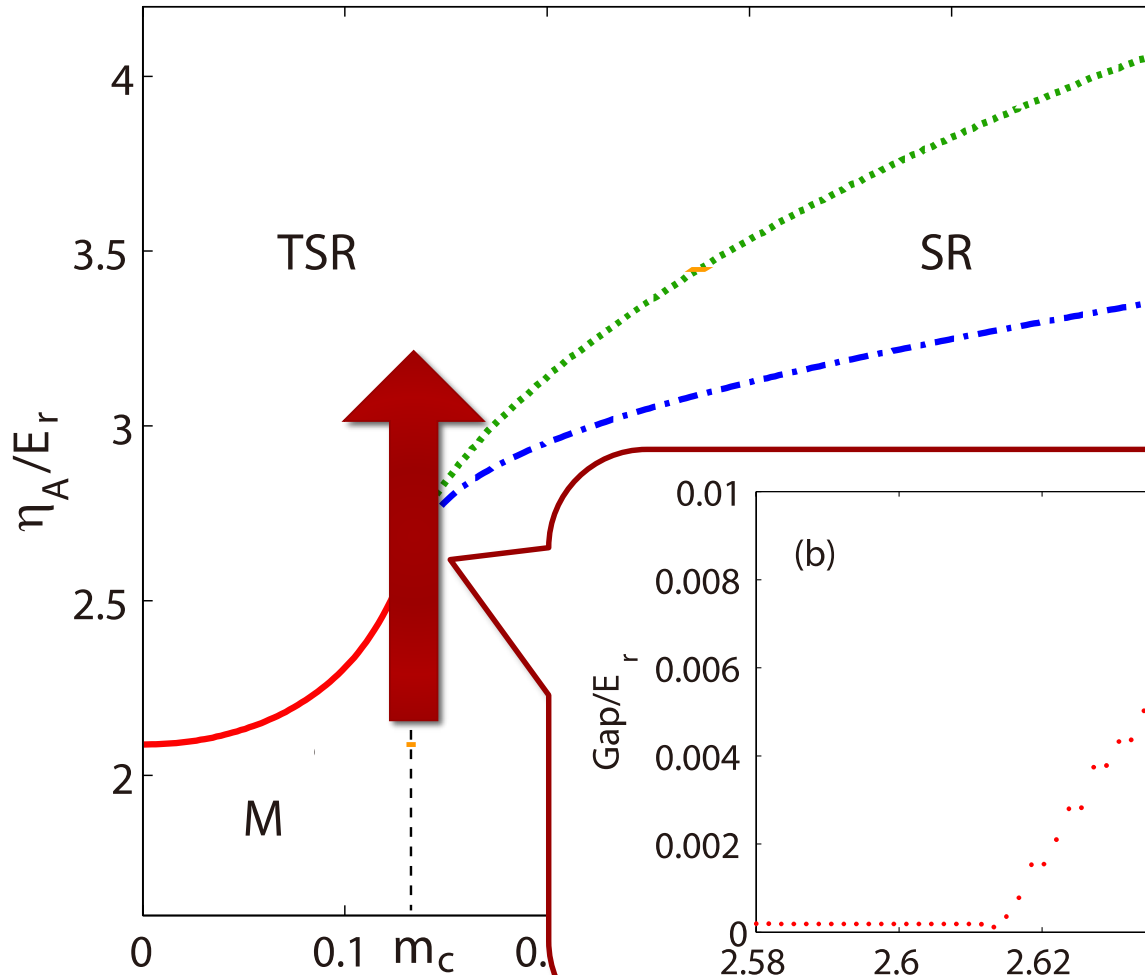


Effective Zeeman Splitting



Tetracritical point

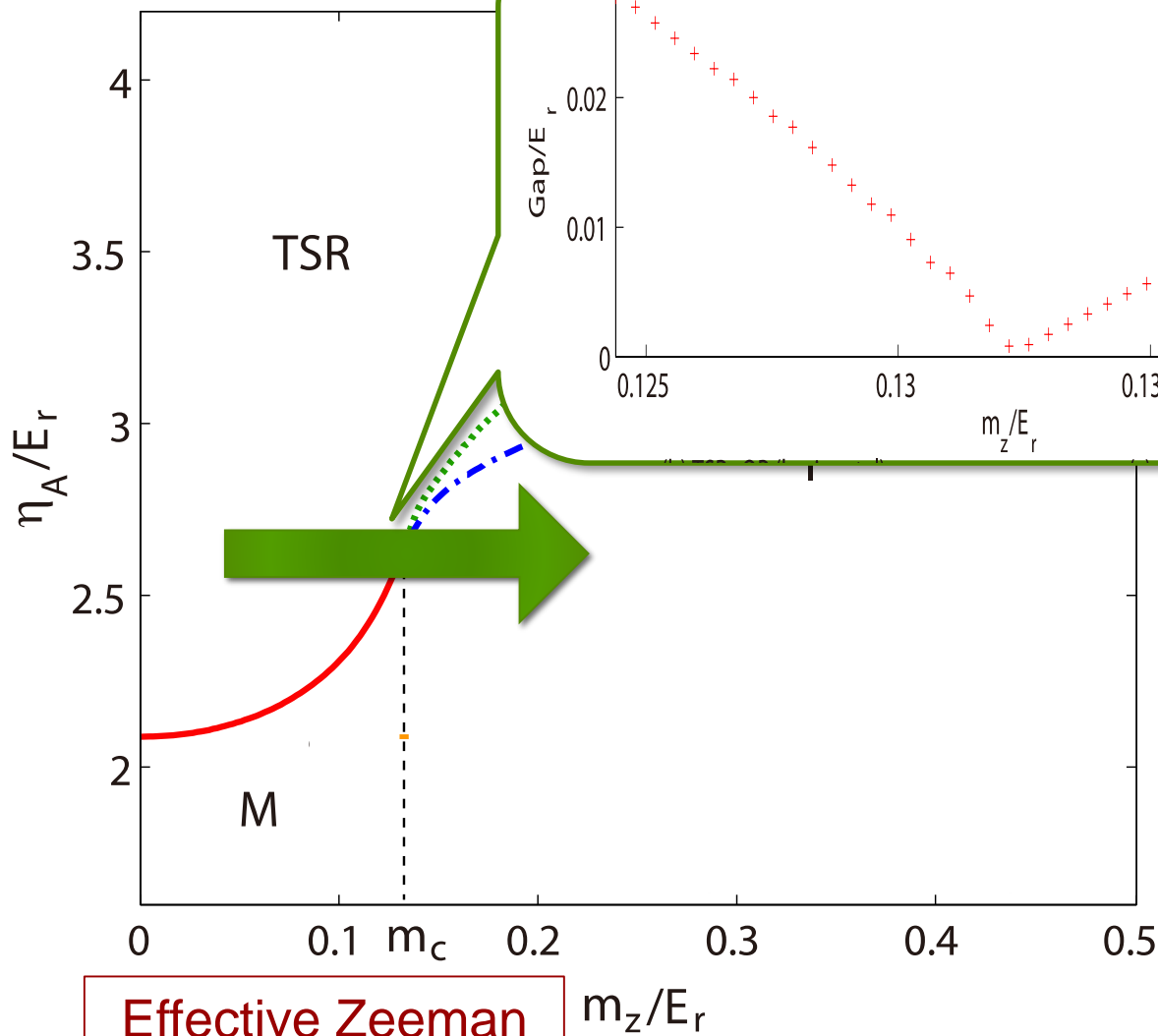
Transverse Pumping



Effective Zeeman Splitting

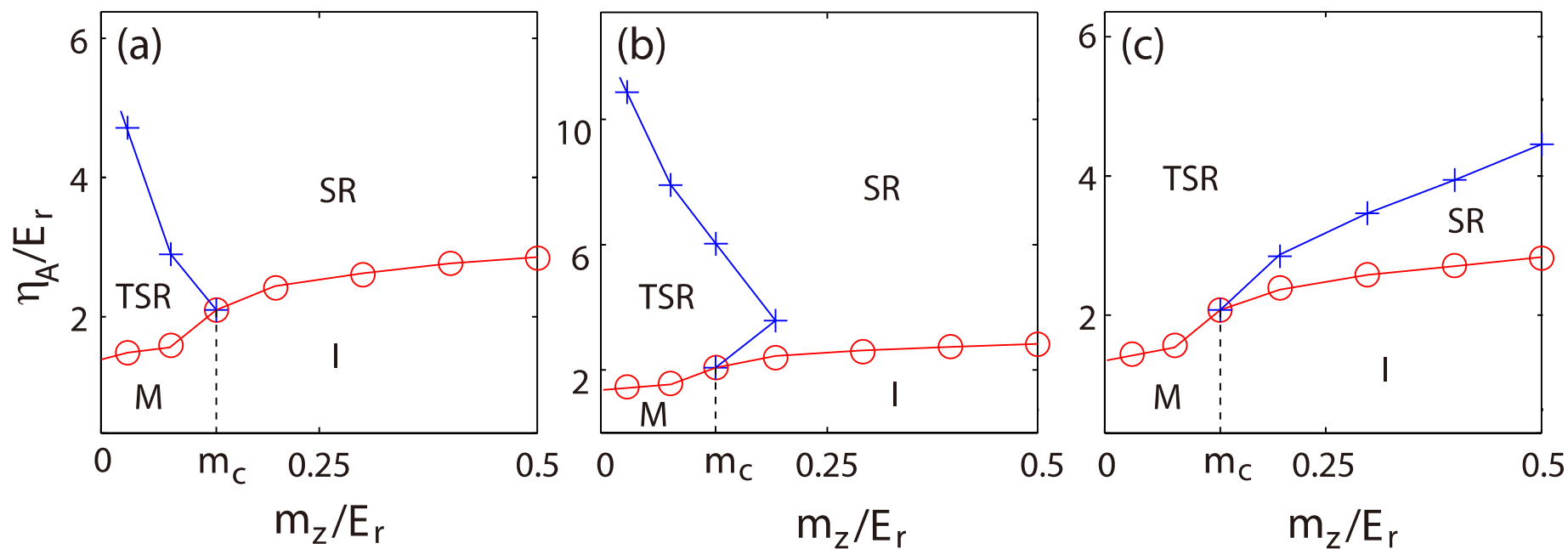
Tetracritical point

Transverse Pumping



Effective Zeeman Splitting

Higher bands effects



1-band approximation

2-band approximation

3-band approximation

Summary

- Superradiance of a quasi-1D non-interacting Fermi gas in a cavity
- Synthetic SOC realized by Raman scheme with driving and cavity laser fields
- Topological SR state

PRL 115, 045303 (2015)
Prog. Phys. 37(4), 125 (2017)
Front. Phys. 13, 136701 (2018)

