

Analysis and numerical approximation of space-fractional boundary value problems in one and two space dimensions

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- Conservation of mass $\partial u/\partial t + \partial q/\partial x = f$ + Fick law $q = -\partial u/\partial x$
 $\implies \partial u/\partial t - \partial^2 u/\partial x^2 = f \implies$ Steady state $-D^2 u(x) = f(x)$.
- Nonlocal Fick law on spatial domain $(0, 1)$: $q = -I_r^{2-\alpha} \partial u/\partial x$ where

$$I_r^{2-\alpha} := r {}_0I_x^{2-\alpha} + (1-r) {}_xI_1^{2-\alpha}, \quad 1 < \alpha < 2, \quad 0 \leq r \leq 1$$

and

$${}_0I_x^{2-\alpha} g := \frac{1}{\Gamma(2-\alpha)} \int_0^x \frac{g(s)}{(x-s)^{\alpha-1}} ds, \quad {}_xI_1^{2-\alpha} g = \frac{1}{\Gamma(2-\alpha)} \int_x^1 \frac{g(s)}{(s-x)^{\alpha-1}} ds$$

$$\implies -DI_r^{2-\alpha} Du(x) = f(x).$$

- When $r = 1/2$, $-DI_r^{2-\alpha} D$ corresponds to $(-\Delta)^{\alpha/2}$.
- Variable-coefficient problems:

$$q = -K(x) I_r^{2-\alpha} \partial u/\partial x \implies -D[K(x) I_r^{2-\alpha} D]u(x) = f(x),$$

$$q = -I_r^{2-\alpha} (K(x) \partial u/\partial x) \implies -DI_r^{2-\alpha} (K(x) Du(x)) = f(x).$$

- Boundary conditions: $u(0) = u(1) = 0$.

- Let $\omega^{(a,b)} = \omega^{(a,b)}(x) := (1-x)^a x^b$ for $a, b > -1$, $\{G_n^{a,b}(x)\}$ be Jacobi polynomials on $(0, 1)$ and $\hat{G}_n^{a,b}(x) = G_n^{a,b}(x) / \|G_n^{a,b}(x)\|$. Let $L_\omega^2(0, 1) := \{g(x) : \int_0^1 \omega(x) g(x)^2 dx < \infty\}$, for weight function $\omega(x) (> 0$ on $(0, 1))$ associated with the inner product $(f, g)_\omega := \int_0^1 \omega(x) f(x) g(x) dx$ and the norm $\|g\|_\omega := (g, g)_\omega^{1/2}$.
- Weighted space $H_{\omega^{(a,b)}}^r := \{v \mid v \text{ is measurable and } \|v\|_{r, \omega^{(a,b)}} < \infty\}$ for $r \in \mathbb{N}$ with associated norm and semi-norm

$$\|v\|_{r, \omega^{(a,b)}} := \left(\sum_{j=0}^r \|D^j v\|_{\omega^{(a+j, b+j)}}^2 \right)^{1/2}, \quad |v|_{r, \omega^{(a,b)}} := \|D^r v\|_{\omega^{(a+r, b+r)}}.$$

$H_{\omega^{(a,b)}}^r$ with non-integer r could be defined by K-method of interpolation.

- Equivalent space: For $r \in \mathbb{R}$, $H_{\omega^{(a,b)}}^r := \{v \in L_{\omega^{(a,b)}}^2 \mid \|v\|_{r, \omega^{(a,b)}} < \infty\}$ where $\|v\|_{r, \omega^{(a,b)}}^2 = \sum_{j=0}^{\infty} (1+j^2)^r v_j^2$ and $\{v_j\}$ are Fourier coefficients of v under the orthonormal basis $\{\hat{G}_j^{a,b}(x)\}$ of $L_{\omega^{(a,b)}}^2$.

- Coercivity: $(I_r^{2-\alpha} Du, Du) \geq c_0 \|u\|_{H^{\alpha/2}}^2$.
- Eigen-relation: $-DI_r^{2-\alpha} D(\omega^{(\alpha-\beta, \beta)} G_n^{\alpha-\beta, \beta}(x)) = \lambda_n G_n^{\beta, \alpha-\beta}(x)$ where $\alpha - 1 \leq \alpha - \beta, \beta \leq 1$ satisfies $(1-r) \sin(\pi\beta) = r \sin(\pi(\alpha - \beta))$ and

$$\lambda_n = \frac{-\sin(\pi\alpha)}{\sin(\pi(\alpha - \beta)) + \sin(\pi\beta)} \frac{\Gamma(n + 1 + \alpha)}{\Gamma(n + 1)}.$$

- $DI_r^{2-\alpha} \omega^{(\alpha-\beta, \beta)} G_n^{(\alpha-\beta, \beta)}(x) = \lambda_n G_{n+1}^{(\beta-1, \alpha-\beta-1)}(x)$ where

$$\lambda_n = \frac{\sin(\pi\alpha)}{\sin(\pi(\alpha - \beta)) + \sin(\pi\beta)} \frac{\Gamma(n + \alpha)}{n!} \sim -(n + 1)^{\alpha-1}.$$

- $\ker(DI_r^{2-\alpha}) = \text{span}\{k(x) := (1-x)^{\alpha-\beta-1} x^{\beta-1}\}$.
- $DI_r^{2-\alpha}(x k(x)) = \lambda_{-1}$ and $DI_r^{2-\alpha}((1-x) k(x)) = -\lambda_{-1}$ where

$$\lambda_{-1} := -(1-r) \Gamma(\alpha) \frac{\sin(\pi\alpha)}{\sin(\pi(\alpha - \beta))}.$$

- Consider the homogeneous boundary value problem of $-DI_r^{2-\alpha}Dw = f$.
- $f(x) \in L^2_{\omega^{(\beta, \alpha-\beta)}}$ could be expressed as $f(x) = \sum_{i=0}^{\infty} f_i G_i^{(\beta, \alpha-\beta)}(x)$.
- If we assume $w(x) = \omega^{(\alpha-\beta, \beta)}(x) \sum_{i=0}^{\infty} w_i G_i^{(\alpha-\beta, \beta)}(x)$, we apply the relation $-DI_r^{2-\alpha}D(\omega^{(\alpha-\beta, \beta)} G_n^{\alpha-\beta, \beta}(x)) = \lambda_n G_n^{\beta, \alpha-\beta}(x)$ to obtain

$$\begin{aligned} -DI_r^{2-\alpha}Dw &= -DI_r^{2-\alpha}D \sum_{i=0}^{\infty} w_i (\omega^{(\alpha-\beta, \beta)}(x) G_i^{(\alpha-\beta, \beta)}(x)) \\ &= \sum_{i=0}^{\infty} w_i \lambda_i G_i^{(\beta, \alpha-\beta)}(x) = \sum_{i=0}^{\infty} f_i G_i^{(\beta, \alpha-\beta)}(x) \implies w_i = \frac{f_i}{\lambda_i}. \end{aligned}$$

- Spectral approximation of w : N -term truncation $w_N(x)$.
- Error estimate: $\|w - w_N\|_{\omega^{-(\alpha-\beta, \beta)}} \leq QN^{-\alpha} \|f\|_{\omega^{(\beta, \alpha-\beta)}}$.
- **Questions** for $-D(KI_r^{2-\alpha}D)u = f$ and $-DI_r^{2-\alpha}(KDu) = f$:
 - FEM: Loss of coercivity of bilinear forms (Wang-Yang SINUM 13).
 - Spectral method: Loss of eigen-relations.

- The solution to $-DI_r^{2-\alpha}(KDu) = f$ could be expressed in terms of w

$$u(x) = C_1 \int_0^x \frac{\omega^{(\alpha-\beta-1, \beta-1)}(s)}{K(s)} ds + \int_0^x \frac{Dw(s)}{K(s)} ds \quad (1)$$

where C_1 is chosen such that $u(1) = 0$.

- A direct verification

$$-DI_r^{2-\alpha}(KDu) = -DI_r^{2-\alpha}(C_1\omega^{(\alpha-\beta-1, \beta-1)}(x) + Dw(x)) = 0 + f(x).$$

- $f(x) \in L^2_{\omega^{(\beta, \alpha-\beta)}} \implies$ unique solution $u \in L^\infty$ and $\|u\|_{L^\infty} \leq Q\|f\|_{\omega^{(\beta, \alpha-\beta)}}$.
- Approximation u_N of u is given by (1) with w replaced by w_N .
- $f \in H^t_{\omega^{(\beta, \alpha-\beta)}}$ for $t \geq 0$ and $K' \in L^2_{\omega^{(\alpha-\beta, \beta)}} \implies$

$$\|u - u_N\|_{\omega^{(-(\alpha-\beta), -\beta)}} \leq QN^{-\alpha-t}\|f\|_{t, \omega^{(\beta, \alpha-\beta)}},$$

$$\|D(u - u_N)\|_{\omega^{(-(\alpha-\beta)+1, -\beta+1)}} \leq QN^{-(\alpha-1)-t}\|f\|_{t, \omega^{(\beta, \alpha-\beta)}}.$$

- Equivalent form of $-D(KI_r^{2-\alpha}D)u = f$ is

$$-I_r^{2-\alpha}Du = f_1 - Af_2$$

where $A := K(0)[I_r^{2-\alpha}Du(x)]|_{x=0}$, $f_1(x) = K^{-1}(x) \int_0^x f(y)dy$ and $f_2(x) = K^{-1}(x)$.

- Spectral expansion $u(x) = \omega^{(\alpha-\beta,\beta)}(x) \sum_{i=0}^{\infty} u_i G_i^{(\alpha-\beta,\beta)}(x)$.
- A useful relation $I_r^{2-\alpha}D(\omega^{(\alpha-\beta,\beta)}G_n^{(\alpha-\beta,\beta)}(x)) = \lambda_n G_{n+1}^{(\beta-1,\alpha-\beta-1)}(x)$.
- $f \in L_{\omega^{(\beta,\alpha-\beta)}}^2 \implies f_1, f_2 \in L_{\omega^{(\beta-1,\alpha-\beta-1)}}^2 \implies$ unique $u \in L_{\omega^{(-(\alpha-\beta),-\beta)}}^2$ such that

$$\|u\|_{\omega^{(-(\alpha-\beta),-\beta)}} \leq C\|f\|_{\omega^{(\beta,\alpha-\beta)}}$$

and

$$A = (f_1, 1)_{\omega^{(\beta-1,\alpha-\beta-1)}} / (K^{-1}, 1)_{\omega^{(\beta-1,\alpha-\beta-1)}}$$

- $f \in H_{\omega^{(\beta,\alpha-\beta)}}^t$ and $K \in W_{\infty}^{t+1}$ for $t \geq 0 \implies$

$$\|u - u_N\|_{\omega^{(-(\alpha-\beta),-\beta)}} \leq Q\|K\|_{W_{\infty}^{t+1}} N^{-\alpha-t} (\|f\|_{t,\omega^{(\beta,\alpha-\beta)}} + 1).$$

- Bilinear form of $-D(KI_r^{2-\alpha}D)u = f$ is $a(g, \hat{g}) := (KI_r^{2-\alpha}Dg, D\hat{g})$.
- Recall that this bilinear form is not coercive.
- Product rule for $-D(K \cdot (I_r^{2-\alpha}Du))$:

$$-D(I_r^{2-\alpha}D)u - \frac{K'}{K}(I_r^{2-\alpha}D)u = \frac{f}{K}$$

- Weak coercivity (Garding inequality): $\|g\|_{L^2}^2 + a(g, g) \geq Q\|g\|_{H^{\alpha/2}}^2$
 \implies Error estimate: $\|u - u_h\|_{L^2} + h^{\sigma-\alpha/2}\|u - u_h\|_{H^{\alpha/2}} \leq Qh^{2\sigma-\alpha}$ where σ is determined by $\|u - I_h u\|_{H^{\alpha/2}} \leq Q\|u\|_{H^\sigma} h^{\sigma-\alpha/2}$.
- Application to $-DI_r^{2-\alpha}(KDu) = f$ by using $KDu = D(Ku) - K'u$
 $\implies -D(I_r^{2-\alpha}D(Ku)) + DI_r^{2-\alpha}(K'u) = f$
 \implies For $v := Ku$ it holds $-D(I_r^{2-\alpha}Dv) + DI_r^{2-\alpha}(K'v/K) = f$.
- Extended to study high-dimensional problems.

- Regularity of $-D(I_r^{2-\alpha}D)w = f$: $f \in L^2 \implies w \in H^{\min\{\beta, \alpha-\beta\}+1/2-\varepsilon}$ (see e.g. Ervin JDE 21).
- Equivalent form of $-D(KI_r^{2-\alpha}D)u = f$ is

$$-DI_r^{2-\alpha}Du = f'_1 - Af'_2$$

where $A := K(0)[I_r^{2-\alpha}Du(x)]|_{x=0}$, $f_1(x) = K^{-1}(x) \int_0^x f(y)dy$ and $f_2(x) = K^{-1}(x) \implies f \in L^2$ implies $u \in H^{\min\{\beta, \alpha-\beta\}+1/2-\varepsilon}$.

- In fact, the above conclusion also holds for $-D(I_r^{2-\alpha}KD)u = f$. Recall that

$$u(x) = C_1 \int_0^x \frac{\omega^{(\alpha-\beta-1, \beta-1)}(s)}{K(s)} ds + \int_0^x \frac{Dw(s)}{K(s)} ds.$$

\implies Key ingredient: $x^\gamma \in H^{\gamma+1/2-\varepsilon}$ for $0 < \gamma < 1$ by interpolation estimate.

- It was shown that the bilinear form of $-DI_r^{2-\alpha}(KD)$ is not coercive and the accuracy of FEM is relatively low, which motivates the (weighted) Petrov-Galerkin method.
- Let $\omega(x) := (1-x)^{\alpha-\beta}x^\beta$ and $\omega^*(x) := (1-x)^\beta x^{\alpha-\beta}$. Based on the structure of $u(x)$ ($u(x) = \omega(x)\phi(x)$ for some $\phi(x)$), we propose a weak formulation of $-DI_r^{2-\alpha}(KD) = f$ via $B : H_\omega^1 \times H_{\omega^*}^{\alpha-1} \rightarrow \mathbb{R}$

$$B(\phi, \psi) := (KD(\omega\phi), I_{1-r}^{2-\alpha}D(\omega^*\psi)) = (f, \omega^*\psi).$$

Here we used the following adjoint property $(I_r^{2-\alpha}g, \hat{g}) = (g, I_{1-r}^{2-\alpha}\hat{g})$.

- Key difficulty: Prove the inf-sup condition in the Banach-Nečas-Babuška Theorem

$$\sup_{\phi \in H_\omega^1} \frac{B(\phi, \psi)}{\|\psi\|_{H_{\omega^*}^{\alpha-1}}} \geq c_0 \|\phi\|_{H_\omega^1}.$$

- Idea: Given $\phi \in H_\omega^1$, find $\psi \in H_{\omega^*}^{\alpha-1}$ such that

$$D(\omega\phi) \cdot (I_{1-r}^{2-\alpha}D(\omega^*\psi)) \geq 0 \text{ or even } D(\omega\phi) = I_{1-r}^{2-\alpha}D(\omega^*\psi),$$

which allows moving K outside of the inner product by its low bound and then using the orthogonality of Jacobi polynomials.

- For $\phi(x) = \sum_{i=0}^{\infty} \phi_i \hat{G}_i^{(\alpha-\beta, \beta)}(x) \in H_{\omega}^1$, define $\Phi_i := -\phi_i/\mu_i$ where $\{\mu_i\}$ are given by

$$I_{1-r}^{2-\alpha} D(\omega^{(\beta, \alpha-\beta)}(x) G_k^{(\beta, \alpha-\beta)}(x)) = \mu_k G_{k+1}^{(\alpha-\beta-1, \beta-1)}(x), \quad k \geq 0.$$

- $\mu_k \sim -(k+1)^{\alpha-1}$ as $k \rightarrow \infty$.
- $\phi(x)$ could be rewritten as $\phi(x) = -\sum_{i=0}^{\infty} \mu_i \Phi_i \hat{G}_i^{(\alpha-\beta, \beta)}(x)$.
- It is also known that for $n \geq 0$

$$D(\omega^{(\alpha-\beta, \beta)}(x) G_n^{(\alpha-\beta, \beta)}(x)) = -(n+1)\omega^{(\alpha-\beta-1, \beta-1)} G_{n+1}^{(\alpha-\beta-1, \beta-1)}(x).$$

- Define $\psi(x)$ by $\psi(x) = \sum_{j=0}^{\infty} (j+1)\Phi_j \hat{G}_j^{(\beta, \alpha-\beta)}(x)$.
- As $\phi \in H_{\omega}^1$, $\|\phi\|_{H_{\omega}^1}^2 = \sum_{i=0}^{\infty} (1+i^2)^1 (\mu_i \Phi_i)^2 \simeq \sum_{i=0}^{\infty} (1+i^2)^{\alpha} \Phi_i^2 < \infty$, then $\|\psi\|_{H_{\omega^*}^{\alpha-1}}^2 = \sum_{j=0}^{\infty} (1+j^2)^{\alpha-1} ((j+1)\Phi_j)^2 \simeq \sum_{j=0}^{\infty} (1+j^2)^{\alpha} \Phi_j^2 < \infty$, i.e., $\psi \in H_{\omega^*}^{\alpha-1}$ and $\|\psi\|_{H_{\omega^*}^{\alpha-1}} \simeq \|\phi\|_{H_{\omega}^1}$.

We invoke the above relations, $K(x) \geq k_0 > 0$ and

$$1 \leq \frac{\|G_{i+1}^{(\alpha-\beta-1, \beta-1)}\|}{\|G_i^{(\beta, \alpha-\beta)}\|} \leq 2$$

to find

$$\begin{aligned} B(\phi, \psi) &= (K(D\omega\phi(x)), I_{1-r}^{2-\alpha} D(\omega^*\psi(x))) \\ &= \left(K\omega^{(\alpha-\beta-1, \beta-1)} \sum_{i=0}^{\infty} (i+1)\mu_i \Phi_i \frac{\|G_{i+1}^{(\alpha-\beta-1, \beta-1)}\|}{\|G_i^{(\beta, \alpha-\beta)}\|} \hat{G}_{i+1}^{(\alpha-\beta-1, \beta-1)}(x), \right. \\ &\quad \left. \sum_{j=0}^{\infty} (j+1)\mu_j \Phi_j \frac{\|G_{j+1}^{(\alpha-\beta-1, \beta-1)}\|}{\|G_j^{(\beta, \alpha-\beta)}\|} \hat{G}_{j+1}^{(\alpha-\beta-1, \beta-1)}(x) \right) \\ &\geq k_0 \sum_{i=0}^{\infty} (i+1)^2 \mu_i^2 \Phi_i^2 \frac{\|G_{i+1}^{(\alpha-\beta-1, \beta-1)}\|^2}{\|G_i^{(\beta, \alpha-\beta)}\|^2} \simeq k_0 \sum_{i=0}^{\infty} (i+1)^2 \mu_i^2 \Phi_i^2 \\ &\simeq k_0 \sum_{i=0}^{\infty} (1+i^2) \phi_i^2 \simeq k_0 \|\phi\|_{H_\omega^1}^2 \simeq k_0 \|\phi\|_{H_\omega^1} \|\psi\|_{H_{\omega^*}^{\alpha-1}}. \end{aligned}$$

- Space-fractional diffusion equation in two space dimension

$$\begin{aligned} \nabla \cdot (-\Delta)^{\frac{\alpha-2}{2}} K(x) \nabla \tilde{u}(x) &= f(x), \quad x \in \Omega, \\ \tilde{u}(x) &= 0, \quad x \in \mathbb{R}^2 \setminus \Omega. \end{aligned}$$

- Ω is the unit disk in \mathbb{R}^2 , $0 < \alpha < 2$, $K = \text{diag}(k_1, k_2)$, and $(-\Delta)^{\frac{\alpha-2}{2}}$ is the Riesz Potential operator defined by

$$(-\Delta)^{-\frac{\beta}{2}} g(x) = \frac{1}{\gamma_d(\beta)} \int_{\mathbb{R}^d} \frac{g(x-y)}{|y|^{d-\beta}} dy, \quad 0 < \beta < d$$

with

$$\gamma_d(\beta) := 2^\beta \pi^{d/2} \Gamma(\beta/2) / \Gamma((d-\beta)/2)$$

- Note that this is different from the Riesz space-fractional diffusion equation

$$-k_1 \frac{\partial^\alpha u}{\partial |x_1|^\alpha} - k_2 \frac{\partial^\alpha u}{\partial |x_2|^\alpha} = f(x_1, x_2), \quad (x_1, x_2) \in \Omega.$$

as $(-\Delta)^{\frac{\alpha-2}{2}}$ is not a one-dimensional fractional integral operator.

- The solid harmonic polynomials in \mathbb{R}^d are the polynomials in d variables which satisfy Laplace equation. In \mathbb{R}^2 the solid harmonic polynomials of degree l can be conveniently written in polar coordinates, $((r, \varphi) : 0 \leq r < \infty, 0 \leq \varphi < 2\pi)$, as $\{r^l \cos(l\varphi), r^l \sin(l\varphi)\}$. Define $\omega^\gamma := (1 - r^2)^\gamma$ and

$$V_{l,1}(x) := r^l \cos(l\varphi), \quad l \geq 0 \text{ and } V_{l,-1}(x) := r^l \sin(l\varphi), \quad l \geq 1.$$

- In \mathbb{R}^2 an orthogonal basis for $L_{\omega^\gamma}^2(\Omega)$ is given via $V_{l,\mu}$ and the Jacobi polynomials $\{P_n^{(\gamma,l)}(\cdot)\}$ on $(-1, 1)$

$$\left\{ \bigcup_{l=0}^{\infty} \bigcup_{n=0}^{\infty} \left\{ V_{l,1}(x) P_n^{(\gamma,l)}(2r^2 - 1) \right\} \right\} \cup \left\{ \bigcup_{l=1}^{\infty} \bigcup_{n=0}^{\infty} \left\{ V_{l,-1}(x) P_n^{(\gamma,l)}(2r^2 - 1) \right\} \right\}.$$
- We assume the solution $\tilde{u}(x) = \omega^{\frac{\alpha}{2}} u(x)$ with

$$u(x) = \sum_{l \geq 1, n \geq 0, \mu \in \{1, -1\}} u_{l,n,\mu} V_{l,\mu}(x) P_n^{(\frac{\alpha}{2}, l)}(2r^2 - 1) + \sum_{n \geq 0} \frac{u_{0,n,1}}{2} V_{0,1}(x) P_n^{(\frac{\alpha}{2}, 0)}(2r^2 - 1)$$
 and similarly express $f(x)$ with coefficients $f_{l,n,\mu}$.
- To determine $u_{l,n,\mu}$ by comparing both sides of the equation, we need to compute $\nabla \cdot (-\Delta)^{\frac{\alpha-2}{2}} K(x) \nabla \left((1 - r^2)^{\frac{\alpha}{2}} V_{l,\mu} P_n^{(\frac{\alpha}{2}, l)}(2r^2 - 1) \right)$.

- Compute $\nabla\left(\left(1-r^2\right)^{\frac{\alpha}{2}} V_{l,\mu} P_n^{\left(\frac{\alpha}{2}, l\right)}\left(2r^2-1\right)\right)$ by properties of Jacobi polynomials and translation between Cartesian and polar coordinates.
- To act $(-\Delta)^{\frac{\alpha-2}{2}}$ on the resulting expression of $K(x)\nabla\left(\left(1-r^2\right)^{\frac{\alpha}{2}} V_{l,\mu} P_n^{\left(\frac{\alpha}{2}, l\right)}\left(2r^2-1\right)\right)$, we employ the following relation.

Theorem (Dyda et al., Constr. Approx., 2017)

For $\delta = d + 2l$, s an integer, $\frac{\alpha}{2} - s > -1$,

$$f(x) = \left(1 - |x|^2\right)_+^{\frac{\alpha}{2}-s} V_{l,\mu}(x) P_n^{\left(\frac{\alpha}{2}-s, \frac{\delta}{2}-1\right)}\left(2|x|^2-1\right),$$

$$\begin{aligned} (-\Delta)^{\frac{\alpha-2}{2}} f(x) &= (-1)^{1-s} 2^{\alpha-2} \frac{\Gamma\left(n+1-s+\frac{\alpha}{2}\right) \Gamma\left(n-1+\frac{\delta+\alpha}{2}\right)}{\Gamma(n+1) \Gamma\left(n+1-s+\frac{\delta}{2}\right)} \\ &\quad \times V_{l,\mu}(x) P_{n+1-s}^{\left(\frac{\alpha}{2}-2+s, \frac{\delta}{2}-1\right)}\left(2|x|^2-1\right). \end{aligned}$$

- We finally differentiate the resulting equation to obtain the expression of $\nabla \cdot (-\Delta)^{\frac{\alpha-2}{2}} K(x)\nabla\left(\left(1-r^2\right)^{\frac{\alpha}{2}} V_{l,\mu} P_n^{\left(\frac{\alpha}{2}, l\right)}\left(2r^2-1\right)\right)$.

By comparing the coefficients on both sides of the resulting equation and rescaling them by

$$d_{l,n} = \frac{\Gamma(n+1+\frac{\alpha}{2})}{\Gamma(n+1)} u_{l,n,1}, \quad \text{and} \quad \tilde{f}_{l,n} = 2^{-(\alpha-2)} \frac{\Gamma(n+1+l)}{\Gamma(n+1+\frac{\alpha}{2}+l)} f_{l,n,1},$$

we obtain

$$(n=0, l=0) : (k_1 + k_2) d_{0,0} = \tilde{f}_{0,0}$$

$$(n=0, l=1) : (2k_1 + (k_1 + k_2)) d_{1,0} = \tilde{f}_{1,0}$$

$$\{(n=0, l)\}_{l \geq 2} : 2(k_1 + k_2) d_{l,0} + (k_1 - k_2) d_{l-2,1} = \tilde{f}_{l,0}$$

$$\{(n, l=0)\}_{n \geq 1} : (k_1 - k_2) d_{2,n-1} + (k_1 + k_2) d_{0,n} = \tilde{f}_{0,n}$$

$$\{(n, l=1)\}_{n \geq 1} : (k_1 - k_2) d_{3,n-1} + (2k_1 + (k_1 + k_2)) d_{1,n} = \tilde{f}_{1,n}$$

$$\{(n, l)\}_{n \geq 1, l \geq 2} : (k_1 - k_2) d_{l+2,n-1} + 2(k_1 + k_2) d_{l,n} + (k_1 - k_2) d_{l-2,n+1} = \tilde{f}_{l,n}.$$

A similar system is generated for the case $\mu = -1$. Thus we only consider the case $\mu = 1$ as above for representation.

Recall the most general case in the above 6 formulas

$$\{(n, l)\}_{n \geq 1, l \geq 2} : (k_1 - k_2) d_{l+2, n-1} + 2(k_1 + k_2) d_{l, n} + (k_1 - k_2) d_{l-2, n+1} = \tilde{f}_{l, n}.$$

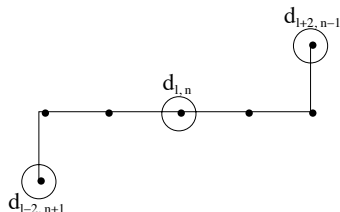


Figure: Stencil illustrating the coupling of the unknowns $d_{l, n}$.

This motivates us to decouple the system into infinite subsystems, each of which is composed by the relations starting from $d_{l+2, n-1} = d_{l_0, 0}, d_{l_0-2, 1}, d_{l_0-4, 2}, \dots$ until $l_0 - 2m = 0$ or $l_0 - 2m = 1$ for some m (Note that since $l_0 - 2m = 0$ and $l_0 - 2m = 1$ correspond to different formulas, thus the coefficient matrices corresponding to even and odd l_0 are slightly different).

- As A_m and $A_{m,*}$ are real, symmetric matrices, and hence have real eigenvalues (singular values), a simple application of Gerschgorin's theorem establishes that there exists constants c_{min} and c_{max} such that the minimum and maximum eigenvalues of the matrices satisfy $0 < c_{min} \leq \lambda_{min}$, $\lambda_{max} \leq c_{max} < \infty$. Hence A_m and $A_{m,*}$ are uniformly invertible with $\|A_m^{-1}\|_2, \|A_{m,*}^{-1}\|_2 \leq c_{min}^{-1}$, i.e. the system is uniquely solvable.
- To estimate the regularity of u , we introduce the weighted function space

$$\mathbf{B}_{\omega^\gamma}^{s_1, s_2}(\Omega) := \left\{ v \mid v \in L_{\omega^\gamma}^2(\Omega) \text{ and } |v|_{\mathbf{B}_{\omega^\gamma}^{s_1, s_2}(\Omega)} < \infty \right\}, \quad s_1, s_2 > 0$$

where the semi-norm $|\cdot|_{\mathbf{B}_{\omega^\gamma}^{s_1, s_2}(\Omega)}$ is defined by

$$|v|_{\mathbf{B}_{\omega^\gamma}^{s_1, s_2}(\Omega)}^2 = \sum_{l, n, \mu} (l^{2s_1} + n^{2s_2}) v_{l, n, \mu}^2 h_{l, n}^2.$$

Here $h_{l, n}$ refers to the norm of the basis function $V_{l, 1}(x) P_n^{(\gamma, l)}(2r^2 - 1)$ and the norm is then defined as

$$\|v\|_{\mathbf{B}_{\omega^\gamma}^{s_1, s_2}(\Omega)}^2 = \sum_{l, n, \mu} (1 + l^{2s_1} + n^{2s_2}) v_{l, n, \mu}^2 h_{l, n}^2.$$

To bound

$$\sum_{l=1}^{\infty} \sum_{n=0}^{\infty} (1 + l^{2s_1} + n^{2s_2}) d_{l,n}^2 h_{l,n}^2$$

we introduce a diagonal matrix W with entries $w_{ii} > 0$ such that

$$W A W^{-1} W \mathbf{d} = W \mathbf{f}, A = A_m \text{ or } A_{m,*}.$$

With $\mathcal{A} = W A W^{-1}$, note that \mathcal{A} is symmetric, positive definite, and has the same eigenvalues as A . Thus, $c_{max}^{-1} \leq \|\mathcal{A}^{-1}\|_2 \leq c_{min}^{-1}$, and $\|W \mathbf{d}\|_2^2 \lesssim \|W \mathbf{f}\|_2^2$, which can be used to for weighted estimates.

Theorem

For $f(x) \in \mathbf{B}_{\frac{\alpha}{2}}^{s_1, s_2}(\Omega)$ there exists a unique solution $\tilde{u}(x) = \omega^{\frac{\alpha}{2}} u(x)$ to the proposed model with $u(x) \in \mathbf{B}_{\frac{\alpha}{2}}^{s_1 + \frac{\alpha}{2}, s_2 + \alpha}(\Omega)$. Furthermore, for $\tilde{g} = \omega^{\frac{\alpha}{2}} g$ with $g \in L_{\omega^{\frac{\alpha}{2}}}^2(\Omega)$, we have

$$(-\Delta)^{\frac{\alpha}{2}} \tilde{g} = -\nabla \cdot (-\Delta)^{\frac{\alpha-2}{2}} \nabla \tilde{g}.$$

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Thank You
for Your Attention!